

4C Vector Calculus

State the value of $\partial x_i / \partial x_j$ and find $\partial r / \partial x_j$, where $r = |\mathbf{x}|$.

[1]
$$\frac{\partial x_i}{\partial x_j} = \delta_{ij}$$

[1]
$$\frac{\partial r}{\partial x_j} = \frac{\partial}{\partial x_j} (x_i x_i)^{1/2} = \frac{1}{2} 2x_j (x_i x_i)^{-1/2} = \frac{x_j}{r}$$

Vector fields \mathbf{u} and \mathbf{v} in \mathbb{R}^3 are given by $\mathbf{u} = r^\alpha \mathbf{x}$ and $\mathbf{v} = \mathbf{k} \times \mathbf{u}$, where α is a constant and \mathbf{k} is a constant vector. Calculate the second-rank tensor $d_{ij} = \partial u_i / \partial x_j$, and deduce that $\nabla \times \mathbf{u} = \mathbf{0}$ and $\nabla \cdot \mathbf{v} = 0$.

[2]
$$\frac{\partial u_i}{\partial x_j} = \frac{\partial}{\partial x_j} (r^\alpha x_i) = \alpha x_j r^{\alpha-2} x_i + r^\alpha \delta_{ij}$$

$$(\nabla \times \mathbf{u})_i = \epsilon_{ijk} \frac{\partial x_k}{\partial x_j} = \epsilon_{ijk} (\alpha x_j x_k r^{\alpha-2} + r^\alpha \delta_{jk}) = 0$$

[2] *as ϵ_{ijk} is antisymmetric in j and k while the term in brackets is symmetric.*

$$\nabla \cdot \mathbf{v} = \frac{\partial}{\partial x_j} (\epsilon_{jkl} k_k u_l) = \epsilon_{jkl} k_k (\alpha r^{\alpha-2} x_l x_j + r^\alpha \delta_{jl}) = 0$$

[2] *as ϵ_{jkl} is antisymmetric in j and l while the term in brackets is symmetric.*

When $\alpha = -3$, show that $\nabla \cdot \mathbf{u} = 0$ and

$$\nabla \times \mathbf{v} = \frac{3(\mathbf{k} \cdot \mathbf{x})\mathbf{x} - \mathbf{k}r^2}{r^5} .$$

$$\nabla \cdot \mathbf{u} = \alpha r^\alpha + r^\alpha \delta_{ii} = -3r^{-3} + 3r^{-3} = 0 \text{ if } \alpha = -3$$

$$\begin{aligned} (\nabla \times \mathbf{v})_i &= \epsilon_{ijk} \frac{\partial}{\partial x_j} (\epsilon_{klm} k_l u_m) \\ &= (\delta_{il} \delta_{jm} - \delta_{im} \delta_{jl}) k_l (-3r^{-5} x_m x_j + r^{-3} \delta_{mj}) \\ &= -3k_i r^{-5} r^2 + 3r^{-5} k_j x_i x_j + 3r^{-3} k_i - k_j r^{-3} \delta_{ij} \\ &= (3r^{-5} (\mathbf{k} \cdot \mathbf{x})\mathbf{x} - \mathbf{k}r^{-3})_i \end{aligned}$$

[2]

[10]

9C Vector Calculus

Write down the most general isotropic tensors of rank 2 and 3. Use the tensor transformation law to show that they are, indeed, isotropic.

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- [2] $\lambda\delta_{ij}$ and $\mu\epsilon_{ijk}$.
 Check that $R_{li}R_{mj}\delta_{ij} = R_{li}R_{mi} = (R^T R)_{lm} = \delta_{mn}$
- [4] and $R_{li}R_{mj}R_{nk}\epsilon_{ijk} = \epsilon_{lmn} \det R = \epsilon_{lmn}$
-

Let V be the sphere $0 \leq r \leq a$. Explain briefly why

$$T_{i_1 \dots i_n} = \int_V x_{i_1} \dots x_{i_n} dV$$

is an isotropic tensor for any n .

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- [2] *A rotation of the basis is equivalent to a backward rotation of the sphere. Spheres have no preferred direction, so the integral will still be T .*
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Hence show that

$$\int_V x_i x_j dV = \alpha \delta_{ij}, \quad \int_V x_i x_j x_k dV = 0 \quad \text{and} \quad \int_V x_i x_j x_k x_l dV = \beta(\delta_{ij}\delta_{kl} + \delta_{ik}\delta_{jl} + \delta_{il}\delta_{jk})$$

for some scalars α and β , which should be determined using suitable contractions of the indices or otherwise.

[You may assume that the most general isotropic tensor of rank 4 is

$$\lambda\delta_{ij}\delta_{kl} + \mu\delta_{ik}\delta_{jl} + \nu\delta_{il}\delta_{jk},$$

where λ , μ and ν are scalars.]

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- [3] *These are all isotropic and symmetric in the indicies. So $\int_V x_i x_j dV$ must be $\alpha \delta_{ij}$. Contract $i = j$: $3\alpha = \int_V x_i x_i dV = \int_V r^4 \sin \theta dr d\theta d\phi = 4\pi a^5/5$ so $\alpha = 4\pi a^5/15$*

- [1] *By isotropy, $\int_V x_i x_j x_k dV = \mu \epsilon_{ijk}$. But the RHS is antisymmetric, so $\mu = 0$ and $\int_V x_i x_j x_k dV = 0$.*

- [2] *By isotropy and hint, the integral must be $\lambda\delta_{ij}\delta_{kl} + \mu\delta_{ik}\delta_{jl} + \nu\delta_{il}\delta_{jk}$. Check symmetry $i \leftrightarrow j$ swaps $\mu \leftrightarrow \nu$ so must have $\mu = \nu$ and similarly $\mu = \lambda$. So $\int_V x_i x_j x_k x_l dV = \beta(\delta_{ij}\delta_{kl} + \delta_{ik}\delta_{jl} + \delta_{il}\delta_{jk})$ for some β .*

Contract $i = j$ and $k = l$: $x_i x_i x_k x_k = r^4$

$$\int_V r^4 r^2 \sin \theta dr d\theta d\phi = \frac{4\pi a^7}{7} \quad \text{and} \quad \beta(\delta_{ii}\delta_{kk} + \delta_{ik}\delta_{ik} + \delta_{ik}\delta_{ik}) = \beta(9 + 3 + 3)$$

- [4] *so $\beta = 4\pi a^7/105$.*
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Deduce the value of

$$\int_V \mathbf{x} \times (\boldsymbol{\Omega} \times \mathbf{x}) dV,$$

where $\boldsymbol{\Omega}$ is a constant vector.

$$\begin{aligned} \int_V \epsilon_{ijk} x_j \epsilon_{klm} \Omega_l x_m dV &= (\delta_{il} \delta_{jm} - \delta_{im} \delta_{jl}) \Omega_l \left(\frac{4\pi a^5}{15} \delta_{jm} \right) \\ &= (3\delta_{il} - \delta_{il}) \Omega_l \left(\frac{4\pi a^5}{15} \right) = \frac{\Omega_l 8\pi a^5}{15} \end{aligned}$$

[2] $so \int_V \mathbf{x} \times (\boldsymbol{\Omega} \times \mathbf{x}) dV = 8\pi a^5 \boldsymbol{\Omega} / 15$

[20]

11C Vector Calculus

The electric field $\mathbf{E}(\mathbf{x})$ due to a static charge distribution with density $\rho(\mathbf{x})$ satisfies

$$\mathbf{E} = -\nabla\phi, \quad \nabla \cdot \mathbf{E} = \frac{\rho}{\epsilon_0}, \quad (1)$$

where $\phi(\mathbf{x})$ is the corresponding electrostatic potential and ϵ_0 is a constant.

(a) Show that the total charge Q contained within a closed surface S is given by Gauss' law

$$Q = \epsilon_0 \int_S \mathbf{E} \cdot d\mathbf{S}.$$

Assuming spherical symmetry, deduce the electric field and potential due to a point charge q at the origin i.e. for $\rho(\mathbf{x}) = q \delta(\mathbf{x})$.

[2] $Q = \int_V \rho dV = \epsilon_0 \int_V \nabla \cdot \mathbf{E} dV = \epsilon_0 \int_S \mathbf{n} \cdot \mathbf{E} dS$ by divergence theorem.

$\rho(\mathbf{x}) = q\delta(\mathbf{x})$ and $\mathbf{E} = E(r)\mathbf{n}$. Hence by Gauss' law $q = 4\pi r^2 \epsilon_0 E(r)$.

[4] So $E(r) = q/4\pi\epsilon_0 r^2$ and $\phi = q/4\pi\epsilon_0 r$.

(b) Let \mathbf{E}_1 and \mathbf{E}_2 , with potentials ϕ_1 and ϕ_2 respectively, be the solutions to (1) arising from two different charge distributions with densities ρ_1 and ρ_2 . Show that

$$\frac{1}{\epsilon_0} \int_V \phi_1 \rho_2 dV + \int_{\partial V} \phi_1 \nabla \phi_2 \cdot d\mathbf{S} = \frac{1}{\epsilon_0} \int_V \phi_2 \rho_1 dV + \int_{\partial V} \phi_2 \nabla \phi_1 \cdot d\mathbf{S} \quad (2)$$

for any region V with boundary ∂V , where $d\mathbf{S}$ points out of V .

$$\begin{aligned} \nabla \cdot (\phi_2 \nabla \phi_1 - \phi_1 \nabla \phi_2) &= \nabla \phi_2 \cdot \nabla \phi_1 + \phi_2 \nabla^2 \phi_1 - \nabla \phi_1 \cdot \nabla \phi_2 - \phi_1 \nabla^2 \phi_2 \\ &= \phi_2 \nabla^2 \phi_1 - \phi_1 \nabla^2 \phi_2 \text{ and (1) gives } \nabla^2 \phi = -\rho/\epsilon_0. \text{ Hence} \end{aligned}$$

[5] $\int_S (\phi_2 \nabla \phi_1 - \phi_1 \nabla \phi_2) \cdot \mathbf{n} dS = \int_V -\phi_2 \frac{\rho_1}{\epsilon_0} + \phi_1 \frac{\rho_2}{\epsilon_0} dV \Rightarrow \text{answer}$

(c) Suppose that $\rho_1(\mathbf{x}) = 0$ for $|\mathbf{x}| \leq a$ and that $\phi_1(\mathbf{x}) = \Phi$, a constant, on $|\mathbf{x}| = a$. Use the results of (a) and (b) to show that

$$\Phi = \frac{1}{4\pi\epsilon_0} \int_{r>a} \frac{\rho_1(\mathbf{x})}{r} dV.$$

[You may assume that $\phi_1 \rightarrow 0$ as $|\mathbf{x}| \rightarrow \infty$ sufficiently rapidly that any integrals over the 'sphere at infinity' in (2) are zero.]

Take $V = \{r > a\}$ and $\rho_2 = q\delta(\mathbf{x})$, so that $\phi_2 = q/4\pi\epsilon_0 r$. Now part (b) gives

$$\frac{1}{\epsilon_0} \int_V \phi_1 \rho_2 dV + \int_{r=a} \Phi \left(\frac{q}{4\pi\epsilon_0 a^2} \right) dS = \frac{1}{\epsilon_0} \int_V \frac{q}{4\pi\epsilon_0 r} \rho_1 dV + \int_{r=a} \frac{q}{4\pi\epsilon_0 a} \nabla\phi_1 \cdot \mathbf{n} dS$$

But $\rho_2 = 0$ in V , and $\int_{r=a} \nabla\phi_1 \cdot \mathbf{n} dS = Q_1/\epsilon_0 = 0$ by part (a) so

$$\frac{\Phi q}{\epsilon_0} = \frac{q}{4\pi\epsilon_0^2} \int_V \frac{1}{r} \rho_1 dV \quad \Rightarrow \quad \text{answer}$$

[9]

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