

MAT3

**MATHEMATICAL TRIPOS****Part III**Friday 13 June 2025 9:00 am to 11:00 am

---

**PAPER 344****THEORETICAL PHYSICS OF SOFT CONDENSED MATTER****Before you begin please read these instructions carefully**

Candidates have TWO HOURS to complete the written examination.

Attempt no more than **TWO** questions.There are **THREE** questions in total.

The questions carry equal weight.

**STATIONERY REQUIREMENTS**

Cover sheet

Treasury tag

Script paper

Rough paper

**SPECIAL REQUIREMENTS**

None

|                                                                                                                                             |
|---------------------------------------------------------------------------------------------------------------------------------------------|
| <p><b>You may not start to read the questions<br/>printed on the subsequent pages until<br/>instructed to do so by the Invigilator.</b></p> |
|---------------------------------------------------------------------------------------------------------------------------------------------|

1 (a) For a closed system comprising a binary fluid mixture of two species, briefly say what is meant by the compositional order parameter  $\phi(\mathbf{r})$ . In what sense this is ‘conserved’? Write down the Landau Ginzburg free energy  $F[\phi]$ , and explain why a linear term in  $\phi$  has no effect.

(b) Consider a binary fluid in which one of the species (only) can show polar ordering of molecular orientations. For the purposes of this question, the polar order parameter can be viewed as a scalar  $p(\mathbf{r})$ , as would anyway be true in one dimension. The following free energy functional can be used to describe such a system:

$$F[\phi, p] = \int \left\{ \frac{a}{2}\phi^2 + \frac{b}{4}\phi^4 + \frac{\kappa_1}{2}(\nabla\phi)^2 + \frac{c}{2}\phi p^2 + \frac{A}{2}p^2 + \frac{B}{4}p^4 + \frac{\kappa}{2}(\nabla p)^2 \right\} d\mathbf{r}.$$

Show that (i) the coupling term  $\frac{c}{2}\phi p^2$  cannot be ignored despite being linear in  $\phi$ ; (ii) the coupling constant  $c$  can be chosen positive without loss of generality. Also, give physical reasons why (iii) there is no term in  $\phi p$ ; and (iv)  $b$  and  $B$  cannot be negative.

(c) Explain why, at mean field level (uniform  $\phi, p$ ) one may construct for the conserved variable  $\phi$  a free energy density  $F/V = f(\phi)$  by unconstrained minimization over  $p$ . Taking positive  $c$ , show that the result is

$$f(\phi) = \frac{a}{2}\phi^2 + \frac{b}{4}\phi^4 - \frac{C}{2}(\phi - \phi_o)^2\theta(\phi_o - \phi)$$

where  $\theta(x) \equiv \frac{1}{2} \left( \frac{|x|}{x} + 1 \right)$  is the Heaviside function, and  $C$  and  $\phi_o$  are constants, which you should find in terms of the parameters in  $F[\phi, p]$ . Interpret  $\phi_o$ .

(d) Restricting attention to parameter choices with  $a > 0$  and for which  $\phi_o = 0$ :

- (i) Sketch  $f(\phi)$  and show that the system will phase separate for  $C > C_c$ , giving the value of  $C_c$ .
- (ii) Explain a graphical construction on  $f(\phi)$  that determines the two binodal densities  $\phi_1, \phi_2$ . Alternatively and for equal credit, construct but do not solve two simultaneous equations for  $\phi_1, \phi_2$  by equating chemical potential and pressure in the two phases. Is this mean-field phase transition continuous or discontinuous?

(e) Now consider  $F[\phi, p]$  for case  $b = \kappa_1 = 0$ . Here the  $\phi$  field can be eliminated by Gaussian integration to give a free energy function  $F[p]$  that governs the statistics of  $p(\mathbf{r})$ .

- (i) Assuming that the result of the Gaussian integration can also be found by setting  $\frac{\delta F[\phi, p]}{\delta \phi(\mathbf{r})} = 0$ , show that

$$F[p] = \int \left\{ \frac{A}{2}p^2 + \frac{\tilde{B}}{4}p^4 + \frac{\kappa}{2}(\nabla p)^2 \right\} d\mathbf{r}$$

with  $\tilde{B} = B - \frac{c^2}{2a}$ . Further assuming that  $A > 0$  and that  $\tilde{B}$  is negligible, identify the functional form of the correlator  $S_p(q) \equiv \langle p_{\mathbf{q}} p_{-\mathbf{q}} \rangle$  at Fourier wavevector  $\mathbf{q}$ .

- (ii) Denoting the Fourier transform of  $S_p(q)$  by  $C_p(r)$ , say what you can about the likely functional form of the real-space correlator for the  $\phi$  field,  $C_\phi(r) \equiv \langle \phi(\mathbf{0})\phi(\mathbf{r}) \rangle - \langle \phi \rangle^2$ . (You are not asked to calculate  $C_p(r)$  or  $C_\phi(r)$  explicitly.)

**2** (a) Starting from a suitable Landau Ginzburg free energy  $F = \int \{f(\phi) + \kappa(\nabla\phi)^2/2\} d\mathbf{r}$ , and stating any assumptions you make, derive the hydrodynamic (noiseless) equations of Model B in the form

$$\begin{aligned}\dot{\phi} &= -\nabla \cdot \mathbf{J}, \\ \mathbf{J} &= -M\nabla(a\phi + b\phi^3 - \kappa\nabla^2\phi).\end{aligned}$$

(b) Show by considering force balance on a 3D spherical droplet of (large) radius  $R$  with positive  $\phi$ , in coexistence with a surrounding phase of negative  $\phi$ , that the coexistence condition is  $\phi = \pm\phi_B + \delta$  where  $\phi_B = (-a/b)^{1/2}$ , and  $\delta = \lambda/R$  where  $\lambda = \frac{\sigma}{\alpha\phi_B}$  with  $\alpha = f''(\phi_B)$  and  $\sigma$  the interfacial tension.

(c) Briefly explain why, if the bulk phase has  $\phi = -\phi_B$  far from the droplet (rather than  $\phi = -\phi_B + \delta$ ), the droplet will slowly evaporate. Show that, writing  $\phi = -\phi_B + g(\mathbf{r})$  in the exterior region and taking a quasi-static approximation, one has  $\nabla^2 g = 0$  with boundary conditions  $g = \delta$  at the droplet surface and  $g = 0$  at infinity. Show also that the current normal to the droplet surface is  $J_n = -M\alpha\nabla_n g$ , with  $\nabla_n$  a derivative in the normal direction. Show further that the normal velocity of the interface obeys  $v_n = -\frac{J_n}{2\phi_B}$ .

(d) Now consider a flat interface in the horizontal  $(x, y)$  plane between bulk coexisting phases arranged so that  $\phi \rightarrow \mp\phi_B$  for  $z \rightarrow \pm\infty$ , which is perturbed by a small height perturbation  $h(x, y)$  in the  $z$  direction. Assume that the same quasi-static approximation holds with boundary condition  $g(x, y, 0) = -\frac{\sigma K}{2\alpha\phi_B}$ , where  $K(x, y)$  is the interfacial mean curvature. Noting that to leading order in small  $h$ ,  $K = (\partial_x^2 + \partial_y^2)h$ , solve  $\nabla^2 g = 0$  in the upper half space for a sinusoidal height perturbation  $h = \eta \sin(qx)$ . Hence find the normal current  $J_n(x, y, z)$  at  $z = 0^+$ .

(e) Supposing this current to be the only contribution to  $v_n = \dot{h}$ , show that  $\dot{\eta} = -\nu|q|^3\eta$  and find the constant  $\nu$ . Confirm that  $\nu > 0$ , give a brief physical argument for its sign. The nonanalytic decay rate  $\eta|q|^3$  at low  $q$  is a signature of nonlocal dynamics for  $h(x, y)$ . Explain why this arises here.

(f) In practice the above calculation for  $\nu$  gives *half* the correct value. Why?

**3** For a polar liquid crystal with order parameter  $\mathbf{p}(\mathbf{r}, t)$  the law of advection is

$$\frac{D\mathbf{p}}{Dt} = (\partial_t + \mathbf{v} \cdot \nabla) \mathbf{p} + \boldsymbol{\Omega} \cdot \mathbf{p} - \xi \mathbf{D} \cdot \mathbf{p} \quad (1)$$

where  $\boldsymbol{\Omega}$  and  $\mathbf{D}$  are respectively the antisymmetric and symmetric parts of the velocity gradient tensor  $\nabla_i v_j$ . Here  $\xi$  is a material-dependent parameter.

(a) Without detailed calculations, explain why the coefficient of the  $\boldsymbol{\Omega}$  term must be unity, as written, rather than a second material-dependent parameter.

(b) By considering the advective free energy increment in a small incompressible displacement  $\mathbf{u} = \mathbf{v} \Delta t$ , show that the stress tensor  $\Sigma_{ij}^p(\mathbf{r})$ , caused by the polar order parameter, obeys  $\Sigma_{ij}^p = \Sigma_{ij}^{(1)} + \Sigma_{ij}^{(2)} + \Sigma_{ij}^{(3)}$ , where

$$\begin{aligned} \nabla_i \Sigma_{ij}^{(1)} &= -p_i \nabla_j h_i \\ \Sigma_{ij}^{(2)} &= (p_i h_j - p_j h_i)/2 \\ \Sigma_{ij}^{(3)} &= \xi(p_i h_j + p_j h_i)/2 \end{aligned}$$

with  $\mathbf{h}(\mathbf{r}) \equiv \delta F / \delta \mathbf{p}(\mathbf{r})$  the molecular field.

(c) A certain laboratory device is designed to create an incompressible, uniaxial extensional flow in three dimensions. This is a uniform flow field  $\mathbf{v}(x, y, z)$  for which  $\boldsymbol{\Omega} = \mathbf{0}$  and  $\mathbf{D}$  is diagonal in  $(x, y, z)$  axes with eigenvalues  $D_{xx} = \gamma$ ,  $D_{yy} = D_{zz} = -\gamma/2$ .

- (i) Assuming that  $\mathbf{p}$  remains uniform in space, and is governed by the usual hydrodynamic equation of motion  $D\mathbf{p}/Dt = -\Gamma \mathbf{h}$ , show that the effect of the flow is to shift  $\mathbf{h} \rightarrow \tilde{\mathbf{h}} \equiv \mathbf{h} + \alpha(p_x, -p_y/2, -p_z/2)$  and find the constant  $\alpha$ .
- (ii) Taking the free energy density for uniform states to be  $\mathbb{F} = \frac{a}{2} p^2 + \frac{b}{4} p^4$  with  $b > 0$ , find a modified free energy of the form  $\mathbb{F}_{\text{mod}} = \frac{1}{2} p_i A_{ij} p_j + \frac{b}{4} p^4$  describing the system under uniform flow conditions. (You are not asked to go beyond mean field theory.) For each sign of  $\alpha$ , identify from  $\mathbb{F}_{\text{mod}}$  the character of the phase transition whereupon  $\mathbf{p}$  first becomes nonzero on decreasing  $a$  and find the critical value  $a_c$  and  $p(a)$ . In each case include a careful statement of what symmetry gets spontaneously broken at the ordering transition.
- (iii) For the case with  $\alpha < 0$  and in the ordered phase, find the order parameter stress  $\Sigma_{ij}^p$  in the flowing system.

**END OF PAPER**