MAMA/301, NST3AS/301, MAAS/301, NST3PHY/1/QFT, MAPY/1/QFT

### MAT3 MATHEMATICAL TRIPOS Part III

Thursday 5 June 2025  $\,$  9:00 am to 12:00 pm  $\,$ 

## **PAPER 301**

# QUANTUM FIELD THEORY

### Before you begin please read these instructions carefully

Candidates have THREE HOURS to complete the written examination.

Attempt no more than **THREE** questions. There are **FOUR** questions in total. The questions carry equal weight.

#### STATIONERY REQUIREMENTS

Cover sheet Treasury tag Script paper Rough paper

#### SPECIAL REQUIREMENTS None

You may not start to read the questions printed on the subsequent pages until instructed to do so by the Invigilator.

$$\mathcal{L} = i\bar{\psi}\not\!\!\!D\psi - \frac{1}{4}F_{\mu\nu}F^{\mu\nu} - i\lambda\bar{\psi}\left[\gamma^{\mu},\gamma^{\nu}\right]F_{\mu\nu}\psi\,.$$

- (a) Are the coupling constants  $\lambda$  and e relevant, marginal or irrelevant? Explain your answer.
- (b) State how  $\psi(x)$ ,  $\bar{\psi}(x)$ , and  $A_{\mu}(x)$  transform under a Lorentz transformation. Briefly explain why the Lagrangian is Lorentz invariant. Show that  $\mathcal{L}$  is real up to total derivatives.
- (c) Derive the equations of motion of the fields.
- (d) Show that the theory is invariant under a local and global U(1) symmetry. Derive the appropriate Noether current  $j^{\mu}$  and charge Q. Show explicitly that Q is conserved on-shell. Compare  $j^{\mu}$  with the source appearing in Maxwell's equation, and comment on any terms they differ by.
- (e) Consider the transformation  $\psi \to e^{i\alpha\gamma_5}\psi$  with  $\alpha$  a constant and  $\gamma^5 = i\gamma^0\gamma^1\gamma^2\gamma^3$ . Under what condition is this a symmetry of  $\mathcal{L}$ ? Does the same condition apply if the fermion is massive? Is it possible to make this symmetry local by suitable choices of e and  $\lambda$ ? Explain your answer.

 $\mathbf{2}$ 

(a) Consider a real vector field  $V^{\mu}(x)$  in four-dimensional Minkowski space, with Lagrangian

 $\mathcal{L} = a \left(\partial_{\mu} V_{\nu}\right) \left(\partial^{\mu} V^{\nu}\right) + b \left(\partial_{\mu} V_{\nu}\right) \left(\partial^{\nu} V^{\mu}\right) + c \left(\partial_{\mu} V^{\mu}\right) \left(\partial_{\nu} V^{\nu}\right) + m^{2} V^{\mu} V_{\mu} ,$ 

where a, b, c, and m are real constants.

(i) Show that  $\mathcal{L}$  can be simplified, up to total derivatives, to

$$\mathcal{L} = \tilde{a} \, V_{\nu} \partial_{\mu} \partial^{\mu} V^{\nu} + \tilde{b} \, V^{\nu} \partial_{\nu} \partial_{\mu} V^{\mu} + m^2 V^{\mu} V_{\mu} \,\,, \tag{(\star)}$$

with  $\tilde{a}$ ,  $\tilde{b}$  real constants.

- (ii) Using  $(\star)$ , derive the equations of motion for  $V^{\mu}$ .
- (iii) Using Noether's theorem, construct the energy-momentum tensor associated to this Lagrangian, and give an expression for the Hamiltonian.
- (iv) Consider the case with  $\tilde{b} = 0$ . Show that for any value of  $\tilde{a}$ , the energy of the system is not positive definite.
- (v) Assuming  $\tilde{a} > 0$ , find a relationship between  $\tilde{b}$  and  $\tilde{a}$  such that the energy of the system is positive definite on-shell, up to total derivatives. Show that for this choice the equations of motion of the vector field imply that  $\partial_{\mu}V^{\mu} = 0$ .
- (b) Another way to constrain  $\tilde{a}$  and b to their physical values, is as follows. A general vector field can always be decomposed as

$$V_{\mu} = A_{\mu} + \partial_{\mu}\pi \; ,$$

where  $A_{\mu}(x)$  satisfies  $\partial_{\mu}A^{\mu} = 0$  (transverse condition), and  $\pi(x)$  is a scalar field.

- (i) Write the Lagrangian ( $\star$ ) in terms of  $A_{\mu}$  and  $\pi$ . Derive the equations of motion for  $A_{\mu}$  and  $\pi$ .
- (ii) Show that the propagator of  $\pi$ , up to a numerical factor, is

$$\langle 0|\mathrm{T}\pi(x)\pi(y)|0\rangle = \int \frac{d^4k}{(2\pi)^4} \frac{i}{m^2} \left(\frac{1}{k^2} - \frac{\tilde{c}}{\tilde{c}k^2 - m^2}\right) e^{-ik\cdot(x-y)} ,$$

where  $\tilde{c}$  is a constant that you should specify. [*Hint: You may use without proof the fact that the propagator is a Green's function.*]

Show that the term involving  $\tilde{c}$  vanishes if  $\tilde{a}$  and b are related as in part (a)(v) of this question.

**3** In this question we will consider a four-dimensional theory consisting of a massive scalar field  $\phi(x)$  and a massive Dirac spinor field  $\psi(x)$ , and the effects of two different interactions between them. The free part of the Lagragian is

$$\mathcal{L}_0 = \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - \frac{1}{2} M^2 \phi^2 + \bar{\psi} (i \partial \!\!\!/ - m) \psi \; .$$

Here M and m are the masses of  $\phi(x)$  and  $\psi(x)$ , respectively. We are using the notation  $\partial = \gamma^{\mu} \partial_{\mu}$  and  $\gamma^{\mu}$  are the gamma matrices.

(a) Consider the interacting theory with Lagrangian  $\mathcal{L}_0 + \mathcal{L}_1$ , with

$$\mathcal{L}_1 = -\lambda \phi \bar{\psi} \psi$$

with  $\lambda$  a coupling constant.

- (i) Write down the classical equations of motion. State, without derivation, the three Schwinger-Dyson equations for this theory.
- (ii) Following the LSZ approach, derive the S-matrix for the decay  $\phi \to \psi \bar{\psi}$  by defining asymptotic states and identifying the appropriate correlation function that enters in the S-matrix.
- (iii) To leading order in the coupling, evaluate the connected contribution to the correlation function appearing in  $\phi \to \psi \bar{\psi}$  by using the Schwinger-Dyson equations. With this, evaluate the *S*-matrix element for this process to leading order, and identify the corresponding amplitude  $\mathcal{A}_{s,s'}$ .
- (iv) Compute the spin-summed/averaged squared matrix element,

$$\mathcal{P} := \frac{1}{4} \sum_{s,s'=1}^{2} |\mathcal{A}_{s,s'}|^2 ,$$

where s and s' are the spins of the final particles.

(b) Now consider instead the interacting theory with Lagrangian  $\mathcal{L}_0 + \mathcal{L}_2$ , with

$$\mathcal{L}_2 = -ig\phi\bar{\psi}\gamma^5\psi \;,$$

with g a coupling constant and  $\gamma^5 = i\gamma^0\gamma^1\gamma^2\gamma^3$ .

What is the S-matrix for  $\phi \to \psi \bar{\psi}$ ? You do not need to derive the result, just state the appropriate modification of your answer in part (a)(iii) of this question.

Evaluate the spin-summed/averaged squared matrix element in this case. In an experiment, would you be able to detect the difference between the effect of the two interaction terms  $\mathcal{L}_1$  and  $\mathcal{L}_2$ ?

**Useful identities:** You may use without proof the following identities related to the mode expansion of free fields:

$$\begin{split} \phi(x) &= \int \frac{d^3 p}{(2\pi)^3} \frac{1}{\sqrt{2\tilde{\omega}_{\vec{p}}}} \left[ a_{\vec{p}} e^{-ip \cdot x} + a_{\vec{p}}^{\dagger} e^{ip \cdot x} \right] \;, \\ \psi(x) &= \int \frac{d^3 p}{(2\pi)^3} \frac{1}{\sqrt{2\omega_{\vec{p}}}} \sum_{r=1}^2 \left[ b_{\vec{p}}^r u^r(\vec{p}) e^{-ip \cdot x} + c_{\vec{p}}^{r\dagger} v^r(\vec{p}) e^{ip \cdot x} \right] \;, \end{split}$$

Part III, Paper 301

with  $\tilde{\omega}_{\vec{p}}^2 = M^2 + \vec{p}^2$  and  $\omega_{\vec{p}}^2 = m^2 + \vec{p}^2$ . Some identities for spinors are

$$(\not p - m)u^{s}(\vec{p}) = 0,$$
  
 $(\not p + m)v^{s}(\vec{p}) = 0,$ 

$$\sum_{s=1}^{2} u^{s}(\vec{p}) \bar{u}^{s}(\vec{p}) = \not p + m ,$$
$$\sum_{s=1}^{2} v^{s}(\vec{p}) \bar{v}^{s}(\vec{p}) = \not p - m .$$

$$\begin{split} u^{s\dagger}(\vec{p})u^{s'}(\vec{p}) &= v^{s\dagger}(\vec{p})v^{s'}(\vec{p}) &= 2\omega_{\vec{p}}\,\delta^{ss'}\,,\\ \bar{u}^{s}(\vec{p})u^{s'}(\vec{p}) &= -\bar{v}^{s}(\vec{p})v^{s'}(\vec{p}) &= 2m\,\delta^{ss'}\,,\\ u^{s\dagger}(\vec{p})v^{s'}(-\vec{p}) &= v^{s\dagger}(\vec{p})u^{s'}(-\vec{p}) &= 0\,. \end{split}$$

4 Suppose that we have a massive scalar field in four dimensions with Lagrangian

$$\mathcal{L} = \frac{1}{2} \partial_{\mu} \Phi \partial^{\mu} \Phi - \frac{m^2}{2} \Phi^2 \; .$$

Consider the field redefinition

 $\Phi(x) = \varphi(x) + \lambda \,\varphi(x)^2 \;,$ 

where  $\lambda$  is a real constant. The scattering amplitudes for particles created by  $\Phi$  or  $\varphi$  should be unchanged if we make this field redefinition. To check that a field redefinition does not change physical observables, do the following:

- (a) Work out the Lagrangian in terms of the redefined field  $\varphi$ . Identify the free and interactions terms in the Lagrangian.
- (b) For the Lagrangian in terms of  $\varphi$ , write, without derivation, the corresponding Feynman rules in momentum space for the S-matrix.
- (c) Consider the scattering process  $\varphi \varphi \to \varphi \varphi$ . Draw the connected Feynman diagrams that contribute to the S-matrix to leading order in  $\lambda$ .
- (d) Based on your answers to part (c) and (b), evaluate the connected contribution to  $\varphi \varphi \rightarrow \varphi \varphi$  scattering at leading order in  $\lambda$ . Comment on your final answer.

[*Hint:* In part (b), you may use the fact that if three scalar fields interact via

$$\mathcal{L}_{\rm int} = \lambda \phi_1(\partial_\mu \phi_2)(\partial^\mu \phi_3) ,$$

then the rule for the vertex is

$$\phi_{1} \underbrace{ \begin{array}{c} \phi_{2} \\ \phi_{2} \\ k_{2} \\ k_{3} \end{array}}_{k_{3}} = (i\lambda)(ik_{2})_{\mu}(ik_{3})^{\mu}(2\pi)^{4}\delta^{4}(k_{1}-k_{2}-k_{3}) .$$

Consider carefully how this can be applied when the three scalars are identical.]

#### END OF PAPER

Part III, Paper 301