MATHEMATICAL TRIPOS Part III

Thursday, 9 June, $2022 \quad 9{:}00 \ \mathrm{am}$ to $11{:}00 \ \mathrm{am}$

PAPER 356

STOCHASTIC PROCESSES IN THEORETICAL PHYSICS AND BIOLOGY

Before you begin please read these instructions carefully

Candidates have TWO HOURS to complete the written examination.

Attempt no more than **TWO** questions. There are **THREE** questions in total. The questions carry equal weight.

STATIONERY REQUIREMENTS

SPECIAL REQUIREMENTS None

Cover sheet Treasury tag Script paper Rough paper

> You may not start to read the questions printed on the subsequent pages until instructed to do so by the Invigilator.

1 Consider a discrete-state continuous-time Markov process with rates w_{mn} and occupation probability P(n,t). The entropy for this system is defined to be

$$S(t) = \sum_{n} P(n, t) \log P(n, t).$$

(a) Derive an expression for the entropy production in this system.

Consider a system with two states, labelled by the index n = 1, 2. The dynamics is modelled as a Markov process with rates $w_{12} = \alpha$ and $w_{21} = \beta$.

(b) Write down the evolution equation for the state probability P(n, t).

(c) Show that the solution for the initial condition P(1,0) = p, P(2,0) = 1 - p is

$$P(1,t) = \frac{1}{\alpha + \beta} \left[\beta + r e^{-(\alpha + \beta)t} \right]$$
$$P(2,t) = \frac{1}{\alpha + \beta} \left[\alpha - r e^{-(\alpha + \beta)t} \right]$$

where $r = \alpha p - \beta (1 - p)$.

(d) Compute the relaxational and steady state entropy production for this system.

(e) Hence, or otherwise, show that this system is always in detailed balance as $t \to \infty$.

Now consider a system with three states, labelled by the index n = 1, 2, 3 with rates $w_{12} = \alpha, w_{21} = \beta, w_{23} = \alpha, w_{32} = \beta, w_{31} = \alpha, w_{13} = \beta.$

(f) Write down the evolution equation for the state probability P(n,t). Show that the solution for initial condition P(1,0) = 1, P(2,0) = 0, P(3,0) = 0 is

$$P(1,t) = \frac{1}{3} \left[1 + 2e^{-3\phi t} \cos\left(\sqrt{3}\psi t\right) \right]$$
$$P(2,t) = \frac{1}{3} \left[1 - 2e^{-3\phi t} \cos\left(\sqrt{3}\psi t - \frac{\pi}{3}\right) \right]$$
$$P(3,t) = \frac{1}{3} \left[1 - 2e^{-3\phi t} \cos\left(\sqrt{3}\psi t + \frac{\pi}{3}\right) \right]$$

where $\phi = (\alpha + \beta)/2$ and $\psi = (\alpha - \beta)/2$.

(g) Show that the entropy production for this system is

$$\dot{S} = (\alpha - \beta) \log \frac{\alpha}{\beta} + \sum_{i} \left[\alpha P(i, t) - \beta P(j, t) \right] \log \frac{P(i, t)}{P(j, t)}$$

where $j = i + 1 \mod 3$ and that the entropy flowing out of the system is

$$S_{neq} = (\alpha - \beta) \log \frac{\alpha}{\beta}.$$

(h) Hence, or otherwise, identify the condition for this system to be out of detailed balance.

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$$dV_t = -\left[U'(X_t) + \gamma V_t\right] dt + \sqrt{2\gamma k_B T} dW_t, \qquad dX_t = V_t dt, \tag{1}$$

where k_B is the Boltzmann constant and T the temperature.

a) Explain the physical system this may describe and the significance of the positive constant γ . Write the Fokker–Planck equation satisfied by the probability density f(x, v, t) of the stochastic process (X_t, V_t) . Provide a heuristic argument at the level of the SDE to eliminate V_t in the limit $\gamma \gg 1$ and obtain an equation for X_t alone. Hence argue that, in this limit, the probability density $p(x,t) = \int_{\mathbb{R}} f(x,v,t) dv$ satisfies

$$\frac{\partial p}{\partial t} = \frac{\partial}{\partial x} \left[D \frac{\partial p}{\partial x} + \phi'(x) p \right].$$
(2)

where $D = k_B T / \gamma$ and $\phi(x) = U(x) / \gamma$.

b) Consider the potential $\phi(x) = \sin(x) - F_0 x$ with $F_0 \ge 0$. Show that if $0 \le F_0 < 1$, $\phi(x)$ has one local minimum and one local maximum in each period, at $m_k = m + 2\pi k$ and $M_k = M + 2\pi k, k \in \mathbb{Z}$, respectively. Sketch $\phi(x)$.

Now consider the function

$$\hat{p}(x,t) = \sum_{k=-\infty}^{\infty} p(x+2\pi k,t), \qquad x \in \Omega = (M_0, M_1).$$
 (3)

Show that \hat{p} is a probability density in Ω that satisfies (2) with periodic boundary conditions. Show that the stationary solution is

$$\hat{p}_s(x) = \frac{J_0 u(x)}{D(1 - e^{-2\pi F_0/D})}$$
 with $u(x) = e^{-\phi(x)/D} \int_x^{x+2\pi} e^{\phi(y)/D} dy$,

where J_0 is a constant. What does J_0 physically represent and how is it determined?

c) The mean velocity of the overdamped process $v = \mathbb{E}(dX_t)/dt$ satisfies

$$v(t) = \int_{-\infty}^{\infty} J(x, t) \mathrm{d}x,$$

where J is the probability flux, $J(x,t) = -D\frac{\partial p}{\partial x} - \phi'(x)p$. Determine the steady-state mean velocity v_s and deduce from it the net direction of motion for $F_0 > 0$.

[QUESTION CONTINUES ON THE NEXT PAGE]

d) For $F_0 \in (0, 1)$ and $k_B T \ll 1$, show that

$$v_s \approx 2\pi (k_+ - k_-),$$

where k_{+} and k_{-} are the escape rates from m_{0} to M_{1} and M_{0} respectively, given by

$$k_{+} = \frac{\sqrt{|\phi''(m_{0})\phi''(M_{1})|}}{2\pi} e^{-\Delta\phi_{+}/D}, \qquad k_{-} = \frac{\sqrt{|\phi''(m_{0})\phi''(M_{0})|}}{2\pi} e^{-\Delta\phi_{-}/D},$$

and $\Delta \phi_+ = \phi(M_1) - \phi(m_0)$, $\Delta \phi_- = \phi(M_0) - \phi(m_0)$ are the potential barriers.

Use k_{\pm} to reduce the continuous process X_t to a continuous-time random walk on the local minima of $\phi(x)$, $\{m_k\}_{k\in\mathbb{Z}}$, providing the transition rates W(k|l) to jump from m_l to m_k for $k, l \in \mathbb{Z}$. What is the distribution of the residence time in each state?

e) Consider now the same process in a bounded domain, so that the reduced model contains a finite number of wells w_0, w_1, \ldots, w_K with absorbing boundaries at each end. Write the pseudocode for an exact algorithm to simulate the random walk for one particle.

Suppose now that we have N independent particles performing the same random walk. Discuss how the implementation of the algorithm may change between the cases $N \ll K$ and $N \gg K$.

Hint: you may use without proof the Laplace's formula that states, given $\Phi(x, y) : \Omega \to \mathbb{R}$ of the form $\Phi(x, y) = \phi_1(x) + \phi_2(y)$ with the global maximum at $(x, y) = (m, M) \in \Omega$,

$$\int_{\Omega} e^{\lambda \Phi(x,y)} \mathrm{d}x \mathrm{d}y \sim \frac{e^{\lambda \Phi(m,M)}}{\lambda} \frac{2\pi}{\sqrt{|\partial_x^2 \Phi(m,M)\partial_y^2 \Phi(m,M)|}}, \quad \text{for } \lambda \gg 1.$$

3 A simple model for bacterial chemotaxis in one dimension consists of a particle (the bacterium) at position $X_t \in \mathbb{R}$ and velocity $V_t \in \{-s, s\}$ at time t, with

$$\mathrm{d}X_t = V_t \mathrm{d}t,\tag{1}$$

and V_t evolves according to a Poisson process: the particle at x changes its velocity from s to -s at rate $\lambda^+(x)$, and from -s to s at rate $\lambda^-(x)$. The turning rates are given by

$$\lambda^{\pm}(x) = \lambda_0 \mp bc'(x),$$

where c(x) describes the concentration of a chemoattractant substance (which is fixed in time). Assume that $|c'(x)| < \lambda_0/b$ for all $x \in \mathbb{R}$.

- a) Is the stochastic process (X_t, V_t) discrete or continuous? Is (X_t, V_t) a Markovian process? What about X_t alone?
- b) Let $p^{\pm}(x,t)dx = \mathbb{P}(X_t \in [x, x + dx), V_t = \pm s)$ be the probability that, at time t, the bacterium is at position x, moving with velocity $\pm s$. Write the system of equations satisfied by $p^+(x,t)$ and $p^-(x,t)$. By considering the total probability density $p(x,t) = p^+(x,t) + p^-(x,t)$ and the flux $j(x,t) = s[p^+(x,t) p^-(x,t)]$, obtain a closed equation for p(x,t).
- c) Consider the large λ_0 limit, and find appropriate scalings for s and b such that the total probability density p(x, t) evolves according to

$$\frac{\partial p}{\partial t} = \frac{\partial}{\partial x} \left(D \frac{\partial p}{\partial x} - \chi p c' \right), \tag{2}$$

Give the expressions of D and χ in terms of the parameters s, λ_0, b , and the conditions on c(x) for (2) to have a stationary solution in \mathbb{R} . What stochastic differential equation does X_t satisfy in this limit?

d) Consider a population of N bacteria at positions X_t^i , for i = 1, ..., N, evolving according to

$$dX_t^i = \chi c'(X_t^i) dt - \frac{1}{N} \sum_{j=1, j \neq i}^N u'(X_t^i - X_t^j) dt + \sqrt{2D} dW_t^i,$$
(3)

where W_t^i are N independent Wiener processes and u(x) is a symmetric function, u(-x) = u(x), and monotonically decreasing for x > 0. What is the physical interpretation of u?

Consider the one-particle probability density $p(x_1, t)$, defined as $p(x_1, t)dx = \mathbb{P}(X_t^1 \in [x_1, x_1 + dx))$. Write down the evolution equation for $p(x_1, t)$ in terms of the two-particle probability density $P_2(x_1, x_2, t)$, defined as

$$P_2(x_1, x_2, t) dx_1 dx_2 = \mathbb{P}(X_t^1 \in [x_1, x_1 + dx_1), X_t^2 \in [x_2, x_2 + dx_2)).$$

Derive a closed equation for $p(x_1, t)$ using the mean-field approximation as $N \to \infty$.

END OF PAPER

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