## MATHEMATICAL TRIPOS Part III

Friday, 10 June, 2022  $\quad 1:30~\mathrm{pm}$  to  $3:30~\mathrm{pm}$ 

## **PAPER 354**

# GAUGE/GRAVITY DUALITY

#### Before you begin please read these instructions carefully

Candidates have TWO HOURS to complete the written examination.

Attempt **BOTH** questions. There are **TWO** questions in total. The questions carry equal weight.

### STATIONERY REQUIREMENTS

#### SPECIAL REQUIREMENTS None

Cover sheet Treasury tag Script paper Rough paper

> You may not start to read the questions printed on the subsequent pages until instructed to do so by the Invigilator.

 $\mathbf{1}$ 

(a) Derive the unitary bound for an *antisymmetric* tensor primary operator  $\mathcal{O}^{ij}$  in a CFT of general dimension  $d \ge 4$ . Show that there is a more restrictive unitary bound when d = 3.

[Hint: in Euclidean signature, the nonzero commutators in the conformal algebra are:

$$\begin{split} [D, P_a] &= iP_a, & [D, K_a] = -iK_a, \\ [P_a, M_{bc}] &= i(\delta_{ab}P_c - \delta_{ac}P_b), & [K_a, M_{bc}] = i(\delta_{ab}K_c - \delta_{ac}K_b), \\ [K_a, P_b] &= 2i(\delta_{ab}D - M_{ab}), \\ [M_{ad}, M_{bc}] &= i(\delta_{ab}M_{cd} - \delta_{ac}M_{bd} - \delta_{bd}M_{ca} + \delta_{cd}M_{ba}). \end{split}$$

with D,  $M_{ab}$ ,  $P_a$  and  $K_a$  being the generators of dilations, rotations, translations, and special conformal transformations respectively.]

(b) A noninteracting massless 2-form potential in  $AdS_6$  is described by the following action:

$$I = \frac{1}{12} \int d^6 \mathbf{x} \sqrt{-g} F_{abc} F^{abc},$$

where  $F_{abc} = A_{ab,c} + A_{bc,a} + A_{ca,b}$ , the potential is antisymmetric:  $A_{ab} = -A_{ba}$ , and there is a gauge-symmetry  $A_{ab} \rightarrow A_{ab} + \alpha_{a,b} - \alpha_{b,a}$  (for an arbitrary 1-form  $\alpha_a$ ). Suppose that this field arises in the bulk dual to a large N holographic CFT<sub>5</sub>. (Assume we are in a weak coupling limit where we can neglect gravity and other interactions.) Write down the wave equation for  $A_{ab}$  in the AdS-Poincaré spacetime:

$$ds^2 = \frac{1}{z^2} \left( dz^2 + \eta_{ij} \, dx^i dx^j \right),$$

where i, j = 0...4. Determine the behavior of the field as  $z \to 0$ , and thereby calculate the dimensions  $\Delta$  of the corresponding (single-trace, primary) source  $J_{ij}$  and operator  $\mathcal{O}^{ij}$ in the CFT under dilations. Explain how you know which one is the source, and which one is the operator.

(c) In a  $CFT_d$  in Minkowski spacetime, the momentum-space stress-tensor is:

$$T_{ij}(\mathbf{p}) = \frac{1}{(2\pi)^{d/2}} \int d^4 \mathbf{x} \, e^{i\mathbf{p}\cdot\mathbf{x}} T_{ij}(\mathbf{x}).$$

Using symmetry and other physical principles, show that (in d > 2, and for a certain range of momenta **p** and **q**) its 2 point function takes the form:

$$\langle 0|T_{ij}(\mathbf{q})T_{kl}(\mathbf{p})|0\rangle \propto \delta^d(\mathbf{p}-\mathbf{q})|\mathbf{p}|^{\beta} \left(P_{ik}P_{jl}+P_{il}P_{jk}-\frac{2}{d}P_{ij}P_{kl}\right),$$

where  $P_{ij} = \eta_{ij} + \frac{p_i p_j}{p^2}$ . (Do not concern yourself with the overall multiplicative factor of the 2 point function, which depends on the particular theory.) Determine the value of the power  $\beta$ . Comment on the range of momenta for which this formula is valid, and what happens outside of this range.

Assuming the CFT is holographic, give a bulk interpretation of the intermediate states  $T_{ij}(\mathbf{p}) |0\rangle$ .

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 $\mathbf{2}$ 

The metric of  $AdS_3$  is:

$$ds^{2} = \frac{R_{AdS}^{2}}{(\cos r)^{2}} \left( -dt^{2} + dr^{2} + (\sin r)^{2} d\phi^{2} \right),$$

where  $0 \leq r < \pi/2$ ,  $\phi$  is periodic with  $0 \leq \phi < 2\pi$ , and  $R_{AdS}$  is the AdS radius.

(a) The entanglement entropy on any interval in a  $CFT_2$  is given by

$$\frac{c}{3}\ln(\epsilon^{-1}) + \text{finite.}$$

where  $\epsilon$  is a UV cutoff. Confirm this formula using the holographic entropy formula. You may select any choice of boundary interval. Note that the CFT<sub>2</sub> has central charge  $c = 3R_{\text{AdS}}/2G$  in units where  $\hbar = 1$ .

[*Hint: The following Laurent series may be useful:*  $\frac{1}{\sin x} = \csc x = \frac{1}{x} + \frac{x}{6} + \frac{7x^3}{360} + \dots$ ]

(b) (i) Evaluate the action of general relativity (with Gibbons-Hawking term)

$$I_{GR} = \frac{1}{16\pi G} \int_{\mathcal{M}} d^3 \mathbf{x} \sqrt{-g} \left(R - 2\Lambda\right) - \frac{1}{8\pi G} \int_{\partial \mathcal{M}} d^2 \mathbf{y} \sqrt{-h} K$$

on the submanifold  $\mathcal{M} \subset \mathrm{AdS}_3$  consisting of points with  $r \leq (\pi/2) - \epsilon$ . Here  $h_{ij}$  is the pullback of the metric to the boundary  $\partial \mathcal{M}$ , whose extrinsic curvature is  $K := \frac{1}{2}n^a g_{ij,a}g^{ij}$ ( $n^a$  being the inward-pointing normal vector). (Since the manifold is time-translation invariant, there will be an overall infinite factor of  $\Delta t$  in your answer.)

(ii) Show that the divergent contribution to I can be eliminated by adding a local counterterm  $I_{c.t.}$  to the boundary. (Do not consider the factor of  $\Delta t$  to be a divergence, as with respect to the CFT this is an IR divergence rather than a UV one.) Calculate this counterterm for the boundary  $\partial M$  explicitly, in the limit where  $\epsilon \to 0$ .

(iii) The remaining finite part (which you are not asked to calculate) can be written in the form:

$$\frac{1}{\Delta t}\lim_{\epsilon \to 0} (I + I_{c.t.}) = -\frac{c}{12},$$

where c is the central charge of the CFT<sub>2</sub>. What is the physical significance of the right hand side of this equation, from the perspective of the boundary theory?

### END OF PAPER