MATHEMATICAL TRIPOS Part III

Friday, 18 June, 2021 $\,$ 12:00 pm to 3:00 pm

PAPER 306

STRING THEORY

Before you begin please read these instructions carefully

Candidates have THREE HOURS to complete the written examination.

Attempt no more than **THREE** questions. There are **FOUR** questions in total. The questions carry equal weight.

You may not start to read the questions printed on the subsequent pages until instructed to do so by the Invigilator. The Polyakov action for the closed string in flat spacetime, with metric $\eta_{\mu\nu}$, is given

$$S[X,h] = -\frac{1}{4\pi\alpha'} \int_{\Sigma} d^2\sigma \sqrt{|h|} h^{ab} \eta_{\mu\nu} \partial_a X^{\mu} \partial_b X^{\nu}.$$

- 1. By considering the variation of the action with respect to the worldsheet metric h_{ab} , find a classical expression for the stress tensor T_{ab} . Show that the stress tensor is traceless.
- 2. Write down an expression for the Nambu-Goto action for the closed string. Show that the Nambu-Goto and Polyakov actions are classically equivalent.
- 3. Consider the path integral expression for the correlation function

$$\langle \phi_1 ... \phi_n \rangle = \mathcal{N} \int \mathcal{D} X \phi_1 ... \phi_n e^{iS_h[X]},\tag{1}$$

where ϕ_i are *n* local fields on the worldsheet, \mathcal{N} is a normalization constant and $S_h[X]$ is the Polyakov action with the worldsheet metric fixed to a particular gauge. Assuming that the fields ϕ_i are invariant under Weyl transformations, show that the correlation function (1) is Weyl-invariant if

$$h^{ab} \langle T_{ab} \phi_1 \dots \phi_n \rangle = 0.$$

4. Now consider a change in the background metric of the target space

$$\eta_{\mu\nu} \to \eta_{\mu\nu} + \varepsilon_{\mu\nu},$$

where $\varepsilon_{\mu\nu}$ is a constant symmetric tensor. Show that the change in the correlation function (1) is given, to leading order, by

$$\int_{\Sigma} d^2 \sigma \langle \mathcal{O}(\sigma, \tau) \phi_1 ... \phi_n \rangle,$$

where $\mathcal{O}(\sigma, \tau)$ is a field you should find. Briefly explain why a change in the worldsheet metric h_{ab} should not change the correlation functions, whereas a change in the target space metric can change the correlation functions.

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After gauge fixing, the the left-moving sector of closed string may be described in terms of the embedding fields $X^{\mu}(z)$ and ghost fields b(z) and c(z), where z is the left-moving worldsheet coordinate.

1. The stress tensor for the embedding fields is given by

$$T_X(z) = -\frac{1}{\alpha'} : \partial X^{\mu} \partial X_{\mu}(z) : .$$

Under infinitesimal conformal transformations, given by a holomorphic vector field v(z), $T_X(z)$ transforms as

$$\delta_v T_X(z) = 2\partial v(z)T_X(z) + v(z)\partial T_X(z) + c_X \partial^3 v(z).$$

By considering the OPE of $T_X(z)$ and $T_X(w)$, calculate the value of the constant c_X . [Hint: You do not need to compute the full $T_X(z)T_X(w)$ OPE.]

2. The ghost stress tensor is given by

$$T_{\rm gh}(z) = -: \partial b(z)c(z): -2: b(z)\partial c(z):.$$

Assuming the form of the transformation of the ghost stress tensor is

$$\delta_v T_{\rm gh}(z) = 2\partial v(z)T_{\rm gh}(z) + v(z)\partial T_{\rm gh}(z) + c_{\rm gh}\partial^3 v(z),$$

calculate the value of the constant $c_{\rm gh}$.

3. Defining the total stress tensor to be

$$T_{\rm tot}(z) = T_X(z) + T_{\rm gh}(z),$$

determine how the total stress tensor transforms under infinitesimal conformal transformations. Hence derive the commutator $[\mathcal{L}_m, \mathcal{L}_n]$, where \mathcal{L}_n are the modes of the total stress tensor. What happens when the dimension of spacetime is 26?

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The closed string tachyon has mass $M^2 = -4/\alpha'$ and the integrated and unintegrated closed string tachyon vertex operators are

$$V = \int d^2z : e^{ik\cdot X(z,\bar{z})}:, \qquad U(z,\bar{z}) =: c(z)\bar{c}(\bar{z})e^{ik\cdot X(z,\bar{z})}:$$

The scattering amplitude of n tachyons at tree level may be written as

$$\mathcal{A}_{n} = g_{c}^{n-2} \Big\langle U(z_{1}, \bar{z}_{1}) U(z_{2}, \bar{z}_{2}) U(z_{3}, \bar{z}_{3}) V_{4} \dots V_{n} \Big\rangle,$$

where g_c is the closed string coupling.

- 1. Briefly explain why \mathcal{A}_n is independent of the coordinates (z_i, \bar{z}_i) .
- 2. Show that the form of the closed string scattering amplitude of n tachyons is given by

$$\mathcal{A}_n = \delta^{26} \left(\sum_{i=1}^n k_i^{\mu} \right) \int d^2 z_4 \dots \int d^2 z_n \, \mathcal{F}_n,$$

where \mathcal{F}_n is a function of the z_i and k_i^{μ} which you should find. You may assume

$$\langle c(z_1)\bar{c}(\bar{z}_1)c(z_2)\bar{c}(\bar{z}_2)c(z_3)\bar{c}(\bar{z}_3)\rangle = |z_1 - z_2|^2|z_2 - z_3|^2|z_3 - z_1|^2$$

3. the Mandelstam variables s, t and u are given by

$$s = -(k_1 + k_2)^2$$
, $t = -(k_1 + k_3)^2$, $u = -(k_1 + k_4)^2$.

By taking $z_1 = 0$, $z_2 = 1$, $z_3 = \infty$ and setting $z_4 \equiv z$ show that the four-point amplitude may be written as

$$\mathcal{A}_4 = \delta^{26} \left(\sum_{i=1}^n k_i^{\mu} \right) g_c^2 \int d^2 z \ |z|^{-\alpha' u/2 - 4} |1 - z|^{-\alpha' t/2 - 4}.$$

By choosing different values for the z_i , explain why the amplitude \mathcal{A}_4 is invariant under the exchange of $s \leftrightarrow t$. 1. An open string has endpoints at $\sigma = 0, \pi$. Starting with the gauge-fixed Polyakov action

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$$S[X,h] = -\frac{1}{4\pi\alpha'} \int_{\Sigma} d^2\sigma\, \eta_{\mu\nu}\partial_a X^\mu \partial^a X^\nu,$$

derive the two possible boundary conditions consistent with the equations of motion.

- 2. Find expressions for the embedding fields $X^{\mu}(\sigma, \tau)$ and $X^{I}(\sigma, \tau)$ of the string in the presence of two parallel D-branes, the worldvolumes of which extend in the $\mu = 0, 1, ..., p$ directions. The D-branes are separated by a distance $x_{2}^{I} - x_{1}^{I}$ in the remaining D - p - 1 directions, where I = p + 1, ..., D.
- 3. Show that the masses M of the open strings described by the above configuration are given by

$$M^{2} = T^{2}|x_{2} - x_{1}|^{2} + \frac{1}{\alpha'}(N-1),$$

where T and N should be found. Describe the massless spectrum of the string and the interpretation of the massless states from the perspective of the D-branes worldvolumes. What happens to the gauge symmetry on the worldvolumes of the D-branes when the D-branes coincide?

END OF PAPER