

MAT3, MAMA, NST3AS

**MATHEMATICAL TRIPOS**      **Part III**

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Thursday, 6 June, 2019 9:00 am to 12:00 pm

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**PAPER 331**

**HYDRODYNAMIC STABILITY**

*Attempt no more than **THREE** questions.*

*There are **FOUR** questions in total.*

*The questions carry equal weight.*

***STATIONERY REQUIREMENTS***

*Cover sheet*

*Treasury Tag*

*Script paper*

*Rough paper*

***SPECIAL REQUIREMENTS***

*None*

<p><b>You may not start to read the questions printed on the subsequent pages until instructed to do so by the Invigilator.</b></p>
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1

Consider infinitesimal, two-dimensional perturbations about a parallel shear flow in an inviscid, stratified fluid in a finite depth domain  $z \in [-L, L]$  between impermeable boundaries:

$$\begin{aligned} \mathbf{u} &= \bar{U}(z)\hat{\mathbf{x}} + \mathbf{u}'(x, z, t); U_{\min} \leq \bar{U} \leq U_{\max}, \\ p &= \bar{p}(z) + p'(x, z, t), \\ \rho &= \bar{\rho}(z) + \rho'(x, z, t), \\ [\mathbf{u}', p', \rho'] &= [\hat{\mathbf{u}}(z), \hat{p}(z), \hat{\rho}(z)] \exp[ik(x - ct)]; \hat{\mathbf{u}} = (\hat{u}, \hat{w}), \end{aligned}$$

where the wavenumber  $k$  is assumed real, the phase speed  $c = c_r + ic_i$  may in general be complex, and gravity  $\mathbf{g} = -g\hat{\mathbf{z}}$  acts in the negative  $z$ -direction.

(a) Applying the Boussinesq approximation appropriately, show that the vertical velocity eigenfunction  $\hat{w}$  satisfies the Taylor-Goldstein equation:

$$\left( \frac{d^2}{dz^2} - k^2 \right) \hat{w} - \frac{\hat{w}}{(\bar{U} - c)} \frac{d^2 \bar{U}}{dz^2} + \frac{N^2 \hat{w}}{(\bar{U} - c)^2} = 0; \quad N^2 = -\frac{g}{\rho_0} \frac{d\bar{\rho}}{dz},$$

where  $N$  is the buoyancy frequency and  $\rho_0$  is an appropriate reference density.

(b) Let  $\hat{w} = (\bar{U} - c)^a q$  for arbitrary real  $a$ .

(i) Show that

$$\begin{aligned} \int_{-L}^L (\bar{U} - c)^{2a} [|q_z|^2 + k^2 |q|^2] dz &= \int_{-L}^L [(N^2 + a(a-1)[\bar{U}_z]^2) (\bar{U} - c)^{2a-2}] |q|^2 dz \\ &\quad + \int_{-L}^L [(a-1)\bar{U}_{zz}(\bar{U} - c)^{2a-1}] |q|^2 dz, \end{aligned}$$

where subscripts denote differentiation with respect to  $z$ .

(ii) Using an appropriate choice of  $a$ , show that the flow must be marginally stable if

$$4N^2 > \left( \frac{d\bar{U}}{dz} \right)^2,$$

for all  $z \in [-L, L]$ .

(iii) Using a different choice of  $a$ , show that  $U_{\min} < c_r < U_{\max}$  if the flow is unstable.

## 2

Consider infinitesimal, two-dimensional perturbations about a parallel shear flow in an inviscid, stratified fluid:

$$\begin{aligned} \mathbf{u} &= \bar{U}(z)\hat{\mathbf{x}} + \mathbf{u}'(x, z, t), \\ p &= \bar{p}(z) + p'(x, z, t), \\ \rho &= \bar{\rho}(z) + \rho'(x, z, t), \\ [\mathbf{u}', p', \rho'] &= [\hat{\mathbf{u}}(z), \hat{p}(z), \hat{\rho}(z)] \exp[ik(x - ct)]; \quad \hat{\mathbf{u}} = (\hat{u}, \hat{w}), \end{aligned}$$

where the wavenumber  $k$  is assumed real, the phase speed  $c = c_r + ic_i$  may in general be complex, and gravity  $\mathbf{g} = -g\hat{\mathbf{z}}$  acts in the negative  $z$ -direction.

- (a) Assume that there is a piecewise constant distribution of background density  $\bar{\rho}$ . Also assume that there is either a piecewise constant distribution or a piecewise linear distribution of background velocity  $\bar{U}$ . Show that the appropriate jump conditions at interfaces, where at least one of the density, vorticity or velocity are discontinuous, are given by:

$$\left[ \frac{\hat{w}}{(\bar{U} - c)} \right]_+^+ = 0; \quad \left[ (\bar{U} - c) \frac{d}{dz} \hat{w} - \hat{w} \frac{d}{dz} \bar{U} - \frac{g\bar{\rho}}{\rho_0} \left( \frac{\hat{w}}{(\bar{U} - c)} \right) \right]_+^+ = 0.$$

In the Boussinesq approximation, you are given that the vertical velocity eigenfunction  $\hat{w}$  satisfies the Taylor-Goldstein equation:

$$\left( \frac{d^2}{dz^2} - k^2 \right) \hat{w} - \frac{\hat{w}}{(\bar{U} - c)} \frac{d^2}{dz^2} \bar{U} + \frac{N^2 \hat{w}}{(\bar{U} - c)^2} = 0; \quad N^2 = -\frac{g}{\rho_0} \frac{d\bar{\rho}}{dz},$$

where  $N$  is the buoyancy frequency and  $\rho_0$  is an appropriate reference density.

- (b) Consider a three-layer flow:

$$\bar{U} = \begin{cases} \frac{\Delta U z}{h} \\ \frac{\Delta U z}{h} \\ \frac{\Delta U z}{h} \end{cases}, \quad \bar{\rho} = \begin{cases} \rho_0 - \frac{\Delta \rho}{2} & z > \frac{h}{2}; \\ \rho_0 & |z| < \frac{h}{2}; \\ \rho_0 + \frac{\Delta \rho}{2} & z < -\frac{h}{2}. \end{cases}$$

- (i) Show that  $\tilde{c} = 2c/\Delta U$  satisfies

$$\tilde{c}^4 - \left( 2 + \frac{J}{\alpha} \right) \tilde{c}^2 + \frac{(2\alpha - J)^2 - J^2 e^{-4\alpha}}{4\alpha^2} = 0,$$

where  $\alpha = kh/2$  and  $J = g\Delta\rho h/[\rho_0\Delta U^2]$ .

- (ii) Hence show that the flow is unstable for

$$\frac{2\alpha}{1 + e^{-2\alpha}} < J < \frac{2\alpha}{1 - e^{-2\alpha}}.$$

- (iii) Interpret this instability in terms of a wave resonance in the limit of large wavenumber.

## 3

Consider the equation

$$\frac{d\mathbf{x}}{dt} = \mathbb{L}\mathbf{x} = \begin{bmatrix} \lambda_1 & 0 \\ 1 & \lambda_2 \end{bmatrix} \begin{bmatrix} x_1 \\ x_2 \end{bmatrix}. \quad (\star)$$

The energy  $E(t)$  of the solution at a given time is defined as  $E = x_1(t)^2 + x_2(t)^2$ .

(a) Show by induction that

$$\mathbb{L}^n = \begin{bmatrix} \lambda_1^n & 0 \\ \frac{\lambda_1^n - \lambda_2^n}{\lambda_1 - \lambda_2} & \lambda_2^n \end{bmatrix},$$

and hence that

$$\mathbb{A}(t) = e^{\mathbb{L}t} = \sum_{n=0}^{\infty} \frac{(\mathbb{L}t)^n}{n!} = \begin{bmatrix} e^{\lambda_1 t} & 0 \\ \frac{e^{\lambda_1 t} - e^{\lambda_2 t}}{\lambda_1 - \lambda_2} & e^{\lambda_2 t} \end{bmatrix}.$$

(b) Confirm that  $\mathbb{A}(t)\mathbf{x}(0)$  is the solution to  $(\star)$ .

(c) Demonstrate that  $\mathbb{L}$  is a non-normal matrix. Find its eigenvectors and demonstrate that they are not orthogonal.

(d) For the case  $\lambda_1, \lambda_2 < 0$ , find the set of initial conditions  $(x_1, x_2)$  where the energy grows immediately. Hence find conditions on  $\lambda_1$  and  $\lambda_2$  under which no energy growth is possible.

(e) Find a general expression for the maximum energy growth  $G(t) = E(t)/E(0)$  at some time  $t$ . If  $\lambda_2 > \lambda_1$ , show that

$$\lim_{t \rightarrow \infty} G = (1 + \alpha^2)e^{2\lambda_2 t}, \quad (\dagger)$$

where the constant  $\alpha^2$  is to be found. Find the initial condition to achieve the optimal long-time growth of  $(\dagger)$  and comment on its relationship to the eigenvector of  $\mathbb{L}$  corresponding to the eigenvalue  $\lambda_2$ .

4

Two-dimensional Rayleigh-Darcy convection in a porous layer is described by the equations

$$\nabla^2 \psi = -R \frac{\partial \theta}{\partial x} \quad \& \quad \frac{\partial \theta}{\partial t} + \frac{\partial \psi}{\partial z} \frac{\partial \theta}{\partial x} - \frac{\partial \psi}{\partial x} \frac{\partial \theta}{\partial z} = \nabla^2 \theta$$

with boundary conditions  $\psi = 0$  and  $\theta = -1$  at the top boundary  $z = 1$  and  $\psi = 0$  and  $\theta = 0$  at the lower boundary  $z = 0$  where  $\psi(x, z, t)$  is the streamfunction,  $\theta(x, z, t)$  is the temperature and  $R$  is the Rayleigh number.

- (a) Show that  $\psi_0 = 0$  and  $\theta_0 = -z$  is the basic conductive solution.
- (b) Taking perturbations  $\psi = \psi_0 + \psi'$  and  $\theta = \theta_0 + \theta'$ , linearize the system. Hence show that with normal modes of the form  $(\psi', \theta')(x, z, t) = (\hat{\psi}, \hat{\theta})(z)e^{st+ikx}$  the eigenvalue  $s$  is given by

$$s = \frac{k^2 R}{k^2 + n^2 \pi^2} - (k^2 + n^2 \pi^2) \quad \text{for } n = 1, 2, \dots$$

Find the lowest value of  $R = R_c$  for convection to occur and confirm that the most unstable mode has a wavelength of 2 in the  $x$ -direction.

- (c) Now consider the weakly nonlinear saturation of the convection for  $R = R_c + \varepsilon^2$  where  $\varepsilon \ll 1$ . By introducing the slow time scale  $T$  such that  $T = \varepsilon^2 t$ , show that the equations can be written as

$$\begin{aligned} \nabla^2 \psi' + R_c \frac{\partial \theta'}{\partial x} &= -\varepsilon^2 \frac{\partial \theta'}{\partial x}, \\ \nabla^2 \theta' - \frac{\partial \psi'}{\partial x} &= \varepsilon^2 \frac{\partial \theta'}{\partial T} + \frac{\partial \psi'}{\partial z} \frac{\partial \theta'}{\partial x} - \frac{\partial \psi'}{\partial x} \frac{\partial \theta'}{\partial z}. \end{aligned}$$

- (d) By assuming expansions

$$\begin{aligned} \psi'(x, z, t) &= \varepsilon A(T) \psi_1(x, z) + \varepsilon^2 A(T)^2 \psi_2(x, z) + \varepsilon^3 A(T)^3 \psi_3(x, z) + \dots \\ \theta'(x, z, t) &= \varepsilon A(T) \theta_1(x, z) + \varepsilon^2 A(T)^2 \theta_2(x, z) + \varepsilon^3 A(T)^3 \theta_3(x, z) + \dots \end{aligned}$$

find  $\theta_1$  if  $\psi_1 = \cos \pi x \sin \pi z$ . By considering the problem at  $O(\varepsilon^2)$ , find  $\psi_2$  and  $\theta_2$ .

- (e) Write down the equations which  $\psi_3$  and  $\theta_3$  must solve, and using the fact that

$$\int_0^1 \int_0^2 \psi_1 \left[ \nabla^4 \psi_3 + R_c \frac{\partial^2 \psi_3}{\partial x^2} \right] dx dz = 0,$$

derive the amplitude equation

$$\frac{dA}{dT} = \alpha A - \beta A^3,$$

where the coefficients  $\alpha$  and  $\beta$  need only be expressed in terms of integrals involving  $\psi_1$ ,  $\theta_1$ ,  $\psi_2$  and  $\theta_2$ .

**END OF PAPER**