

#### MATHEMATICAL TRIPOS Part III

Monday, 30 May, 2016 9:00 am to 11:00 am

### **PAPER 318**

#### MAGNETOHYDRODYNAMICS

Attempt **ALL** questions. There are **THREE** questions in total. The questions carry equal weight.

STATIONERY REQUIREMENTS

Cover sheet Treasury Tag Script paper **SPECIAL REQUIREMENTS** None

You may not start to read the questions printed on the subsequent pages until instructed to do so by the Invigilator.

## CAMBRIDGE

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Conducting viscous fluid with density  $\rho$ , kinematic viscosity  $\nu$  and magnetic diffusivity  $\eta$  flows steadily in the *x*-direction in the channel bounded by the planes  $z = \pm L$  driven by a uniform pressure gradient  $dp/dx = -\rho G$ . There is a uniform magnetic field  $B_0 \hat{\mathbf{z}}$  applied across the channel. The boundary conditions at  $z = \pm L$  are that the velocity vanishes and the magnetic field is normal to the boundary.

(i) Justify the assertion that the velocity  $\mathbf{u}$  and induced magnetic field  $\mathbf{b}$  (so that the total field is  $B_0\hat{\mathbf{z}} + \mathbf{b}$ ) are of the form  $\mathbf{u} = (u(z), 0, 0)$ ,  $\mathbf{b} = (b(z), 0, 0)$ , where u, b satisfy the coupled equations

$$0 = G + \frac{1}{\mu_0 \rho} B_0 \frac{db}{dz} + \nu \frac{d^2 u}{dz^2}, \quad 0 = B_0 \frac{du}{dz} + \eta \frac{d^2 b}{dz^2},$$

with  $u(\pm L) = b(\pm L) = 0$ .

(ii) Solve these equations for u(z) and b(z) and obtain the expressions

$$\begin{split} u &= \frac{GL^2}{\nu H^2} (H \coth H) \left(1 - \frac{\cosh(Hz/L)}{\cosh H}\right), \\ b &= \frac{B_0}{\eta} \frac{GL^3}{\nu H^2} \left(-\frac{z}{L} + \frac{\sinh(Hz/L)}{\sinh H}\right), \end{split}$$

where *H* is the Hartmann number  $H \equiv B_0 L/\sqrt{\mu_0 \rho \nu \eta}$ . Derive an expression for the flux of fluid  $Q = \int_{-L}^{L} u \, dz$  and sketch its form as a function of *H* keeping the other parameters fixed. Sketch also u(z) and b(z) when  $H \gg 1$ .

### UNIVERSITY OF

 $\mathbf{2}$ 

A simple model of Parker dynamo waves, with a fluctuating  $\alpha$ -effect and with diffusion ignored, can be written in appropriate units as the pair of complex equations (with D > 0)

$$\frac{dA}{dt} = D^2 f(t)B, \qquad \frac{dB}{dt} = \mathrm{i}A,$$

where  $f(t) = 1, 0 < t < T, 2T < t < 3T, \dots$  and  $f(t) = -1, T < t < 2T, 3T < t < 4T, \dots$ 

Consider the period 0 < t < T. Defining  $\sigma = D(1 + i)/\sqrt{2}$ , show that the general solution is

$$\begin{pmatrix} A \\ B \end{pmatrix} = P_+ \begin{pmatrix} \sigma \\ i \end{pmatrix} e^{\sigma t} + Q_+ \begin{pmatrix} -\sigma \\ i \end{pmatrix} e^{-\sigma t},$$

where  $P_+, Q_+$  are complex constants. Deduce that

$$\begin{pmatrix} A(T) \\ B(T) \end{pmatrix} = \begin{pmatrix} \cosh \sigma T & -\mathrm{i}\sigma \sinh \sigma T \\ \frac{\mathrm{i}}{\sigma} \sinh \sigma T & \cosh \sigma T \end{pmatrix} \begin{pmatrix} A(0) \\ B(0) \end{pmatrix} \equiv \mathsf{M}^+ \begin{pmatrix} A(0) \\ B(0) \end{pmatrix}.$$

It may be assumed analogously that for the period T < t < 2T we have the related results

$$\begin{pmatrix} A \\ B \end{pmatrix} = P_{-} \begin{pmatrix} \sigma^* \\ i \end{pmatrix} e^{\sigma^* t} + Q_{-} \begin{pmatrix} -\sigma^* \\ i \end{pmatrix} e^{-\sigma^* t},$$

where  $\sigma * = -i\sigma = D(1-i)/\sqrt{2}$  is the complex conjugate of  $\sigma$ , and

$$\begin{pmatrix} A(2T) \\ B(2T) \end{pmatrix} = \begin{pmatrix} \cosh \sigma^* T & -\mathrm{i}\sigma^* \sinh \sigma^* T \\ \frac{\mathrm{i}}{\sigma^*} \sinh \sigma^* T & \cosh \sigma^* T \end{pmatrix} \begin{pmatrix} A(T) \\ B(T) \end{pmatrix} \equiv \mathsf{M}^- \begin{pmatrix} A(T) \\ B(T) \end{pmatrix}.$$

Show that the matrices  $M^+$ ,  $M^-$  have determinant 1.

Deduce that the mean exponential growth rate of A, B over very large times is given by  $\ln \Lambda/2T$ , where  $\Lambda$  is the larger eigenvalue of the matrix  $\mathsf{N} = \mathsf{M}^-\mathsf{M}^+$ . Determine the determinant and trace of  $\mathsf{N}$  in terms of D, T and hence find  $\Lambda$ . Setting T = 1 show that as  $D \to \infty$  the growth rate  $\to (\sqrt{2}D - \ln 2)/2$ .

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Consider mean field dynamo action due to a time-dependent solenoidal velocity field

$$\boldsymbol{u} = \mathbb{R} \left( \boldsymbol{u}_1 \cos \Omega t + \boldsymbol{u}_2 \sin \Omega t \right)$$

where  $\boldsymbol{u}_i = \hat{\boldsymbol{u}}_i e^{i\boldsymbol{k}_i \cdot \boldsymbol{x}}$  for distinct vectors  $\boldsymbol{k}_1, \boldsymbol{k}_2$  with  $|\boldsymbol{k}_1| = |\boldsymbol{k}_2| = k$ . Assume that  $\boldsymbol{B} = \boldsymbol{B}_0 + \boldsymbol{b}(\boldsymbol{x}, t)$  where  $\boldsymbol{B}_0$  is a constant vector, and that products of fluctuating quantities can be neglected in the equation for  $\boldsymbol{b}$ . Show that  $\boldsymbol{b}$  takes the form

$$\boldsymbol{b} = \mathbb{R} \sum_{j=1,2} e^{i\boldsymbol{k}_j \cdot \boldsymbol{x}} (\boldsymbol{p}_j \cos \Omega t + \boldsymbol{q}_j \sin \Omega t)$$

and determine the  $p_j$  and  $q_j$ .

Using the result (which should be justified)

$$\frac{1}{2}\mathbb{R}(\mathrm{i}\hat{\boldsymbol{u}}_{j}^{*}\times\hat{\boldsymbol{u}}_{j})=H_{j}\boldsymbol{k}_{j}, \ j=1,2$$

for some constants  $H_j$ , show that if the emf  $\mathcal{E} = \langle u \times b \rangle \equiv \alpha \cdot B_0$ , (where  $\langle \cdot \rangle$  denotes an average over space and time), then  $\alpha$  takes the form

$$\alpha_{ij} = \frac{\eta k^2}{2(\Omega^2 + \eta^2 k^4)} (H_1 k_{1i} k_{1j} + H_2 k_{2i} k_{2j}).$$

Give the form of  $\boldsymbol{\alpha}$  when  $\boldsymbol{k}_1 = (k, 0, 0)$ ,  $\boldsymbol{k}_2 = (0, k, 0)$ ,  $H_1 = H_2$ . Is mean field dynamo dynamo action possible? Give brief reasons for your answer (detailed calculation is not required).

#### END OF PAPER