#### MATHEMATICAL TRIPOS Part III

Thursday, 2 June, 2016  $\,$  1:30 pm to 4:30 pm

#### **PAPER 304**

### ADVANCED QUANTUM FIELD THEORY

Attempt no more than **THREE** questions. There are **FOUR** questions in total. The questions carry equal weight.

#### STATIONERY REQUIREMENTS

Cover sheet Treasury Tag Script paper **SPECIAL REQUIREMENTS** None

You may not start to read the questions printed on the subsequent pages until instructed to do so by the Invigilator.

1

Write down the momentum space Feynman rules for the d-dimensional Euclidean scalar theory with action

 $\mathbf{2}$ 

$$S[\phi] = \int \left(\frac{1}{2}\partial_{\mu}\phi \,\partial^{\mu}\phi + \frac{m^2}{2}\phi^2 + \frac{\lambda}{4!}\phi^4\right) d^dx \,,$$

where  $\lambda$  is a coupling constant. Draw all connected, 1-loop, one particle irreducible Feynman graphs with four external lines.

Show that in *d*-dimensional Euclidean momentum space

$$\int \frac{d^d k}{(2\pi)^d} \frac{1}{(k^2 + m^2)^n} = \frac{1}{(n-1)!} \frac{(m^2)^{d/2-n}}{(4\pi)^{d/2}} \Gamma(n-d/2) \,,$$

where  $k^2 = k_{\mu}k^{\mu}$ . Use this result to verify that

$$\int \frac{d^d k}{(2\pi)^d} \frac{1}{(k^2 + m^2)((p-k)^2 + m^2)} \sim \frac{1}{8\pi^2 \epsilon} \quad \text{as } \epsilon = 4 - d \to 0.$$

Hence or otherwise compute the  $\beta$ -function for the quartic coupling in  $S[\phi]$  to 1-loop accuracy in the  $\overline{\text{MS}}$  scheme, and comment on its sign.

[You may assume that  $\int d^d k \, e^{-ak^2} = (\pi/a)^{d/2}$  for a > 0, and that the  $\Gamma$  function  $\Gamma(z) = \int_0^\infty ds \, s^{z-1} \, e^{-s}$  obeys  $\Gamma(1) = 1$  and  $\Gamma(1+z) = z\Gamma(z)$ .]

 $\mathbf{2}$ 

Classical Yang-Mills theory has action

$$S[A] = \frac{1}{4g^2} \int F^a_{\mu\nu} F^{a\mu\nu} \, d^4x \,,$$

where

$$F^a_{\mu\nu} = \partial_\mu A^a_\nu - \partial_\nu A^a_\mu + f^a_{bc} A^b_\mu A^c_\nu$$

is the field strength tensor for the gauge field  $A^a_{\mu}$ . Obtain the Yang-Mills equations and the Bianchi identity obeyed by  $F^a_{\mu\nu}$ .

Without detailed calculation, briefly explain why these equations give a poor description of the phenomenology of quantum Yang-Mills theory at low energies.

Write down the additional terms that must be added to the above action if we wish to quantize Yang-Mills theory in some gauge  $f^a[A] = 0$ . State clearly the quantum numbers of any fields you introduce, and give a brief explanation of their role. What is the significance of choosing  $f^a[A]$  to be a *local* functional?

In the case  $f^{a}[A] = n^{\mu}A^{a}_{\mu}$  where  $n^{\mu}$  is a fixed, constant vector, obtain the Feynman rules for the additional fields, and the (tree-level) momentum space propagator for the gauge field. Briefly explain why the resulting perturbation theory is Lorentz invariant.

3

Consider the action

$$S(z,\bar{z},\theta_i,\bar{\theta}_i) = |W'(z)|^2 - W''(z)\,\theta_1\theta_2 - \overline{W''(z)}\,\bar{\theta}_1\bar{\theta}_2$$

4

for a d = 0 theory, where W(z) is a polynomial in the bosonic variable z and  $W'(z) = \partial W/\partial z$ . The variables  $\theta_1$ ,  $\theta_2$ ,  $\bar{\theta}_1$  and  $\bar{\theta}_2$  are Grassmann-valued.

Show that the Grassmann vector field

$$V_1 = \theta_1 \frac{\partial}{\partial z} - \overline{W'(z)} \frac{\partial}{\partial \theta_2}$$

obeys  $\{V_1, V_1\} = 0$  and  $V_1(S) = 0$ . Find three further, independent Grassmann vector fields that also leave the action invariant.

Evaluate:

- i) the partition function  $Z = \int dz \, d\bar{z} \, d^2\theta \, d^2\bar{\theta} \, e^{-S(z,\bar{z},\theta_i,\bar{\theta}_i)}$ ,
- ii) the (unnormalized) expectation value  $\langle W'(z)g(z)\rangle$  where g(z) is a polynomial, and
- iii) the (unnormalized) expectation value  $\langle f(z) \rangle$ , where f(z) is a further polynomial. [*Hint: Consider the effect of rescaling* W(z).]

 $\mathbf{4}$ 

Consider the behaviour of a theory defined with UV cutoff  $\Lambda_0$  as it flows down to a scale  $\Lambda < \Lambda_0$ . Write down the Callan–Symanzik equation describing the evolution of a correlation function  $\langle \mathcal{O}_1(x_1) \mathcal{O}_2(x_2) \cdots \mathcal{O}_n(x_n) \rangle$  of composite local operators  $\mathcal{O}_i(x_i)$ , explaining the meaning of each of the terms.

A certain massless, four-dimensional theory has a single dimensionless coupling  $g(\Lambda)$ at scales  $\Lambda < \Lambda_0$ . It also has a conserved current  $J_{\mu}$ . The two–point function  $G_0$  of these currents is defined by

$$(k_{\mu}k_{\nu} - \delta_{\mu\nu}k^2) G_0(k^2, g(\Lambda), \Lambda, \Lambda_0) = \int d^4x \, e^{-ik \cdot x} \langle J_{\mu}(x) J_{\nu}(0) \rangle \, d^4x \, e^{-ik \cdot x} \langle J_{\mu}(x) J_{\nu}(0) \rangle \, d^4x \, e^{-ik \cdot x} \langle J_{\mu}(x) J_{\nu}(0) \rangle \, d^4x \, e^{-ik \cdot x} \langle J_{\mu}(x) J_{\nu}(0) \rangle \, d^4x \, e^{-ik \cdot x} \langle J_{\mu}(x) J_{\nu}(0) \rangle \, d^4x \, e^{-ik \cdot x} \langle J_{\mu}(x) J_{\nu}(0) \rangle \, d^4x \, e^{-ik \cdot x} \langle J_{\mu}(x) J_{\nu}(0) \rangle \, d^4x \, e^{-ik \cdot x} \langle J_{\mu}(x) J_{\nu}(0) \rangle \, d^4x \, e^{-ik \cdot x} \langle J_{\mu}(x) J_{\nu}(0) \rangle \, d^4x \, e^{-ik \cdot x} \langle J_{\mu}(x) J_{\nu}(0) \rangle \, d^4x \, e^{-ik \cdot x} \langle J_{\mu}(x) J_{\nu}(0) \rangle \, d^4x \, e^{-ik \cdot x} \langle J_{\mu}(x) J_{\nu}(0) \rangle \, d^4x \, e^{-ik \cdot x} \langle J_{\mu}(x) J_{\nu}(0) \rangle \, d^4x \, e^{-ik \cdot x} \langle J_{\mu}(x) J_{\nu}(0) \rangle \, d^4x \, e^{-ik \cdot x} \langle J_{\mu}(x) J_{\nu}(0) \rangle \, d^4x \, e^{-ik \cdot x} \langle J_{\mu}(x) J_{\nu}(0) \rangle \, d^4x \, e^{-ik \cdot x} \, d^4x \, d^4x \, e^{-ik \cdot x} \, d^4x \, d^$$

What is the classical mass dimension of  $G_0$ ?

 $G_0$  diverges logarithmically as the cut-off  $\Lambda_0 \to \infty$ . At one-loop, a finite, renormalized two-point function  $G_{\rm ren}(k^2, g(\Lambda), \Lambda)$  may be defined by

$$G_{\rm ren}(k^2, g(\Lambda), \Lambda) = \lim_{\Lambda_0 \to \infty} \left[ G_0(k^2, g(\Lambda), \Lambda, \Lambda_0) - a \ln\left(\frac{\Lambda_0}{\Lambda}\right) \right] \,,$$

where a is a constant. Write down the Callan–Symanzik equation obeyed by  $G_{\text{ren}}(k^2, g(\Lambda), \Lambda)$  assuming that the current  $J_{\mu}(x)$  has vanishing anomalous dimensions.

Determine the leading behaviour of  $G_{\rm ren}$  when  $k^2 \gg \Lambda^2$  in the case that the theory is asymptotically free.

#### END OF PAPER