

MATHEMATICAL TRIPOS Part III

Friday, 30 May, 2014 9:00 am to 12:00 pm

PAPER 44

ADVANCED QUANTUM FIELD THEORY

*Attempt no more than **THREE** questions.*

*There are **FOUR** questions in total.*

The questions carry equal weight.

STATIONERY REQUIREMENTS

Cover sheet

Treasury Tag

Script paper

SPECIAL REQUIREMENTS

None

<p>You may not start to read the questions printed on the subsequent pages until instructed to do so by the Invigilator.</p>

1 Let M be an $N \times N$ Hermitian matrix and let

$$V(M) = \frac{1}{2}\text{tr}(M^2) + \frac{g}{4}\text{tr}(M^4)$$

where g is a constant.

- a) Show that $V(M)$ is invariant under $M \rightarrow U^\dagger M U$ where U is a unitary matrix, and hence that $V(M)$ depends only on the eigenvalues $\{\lambda_i\}$ of M .
- b) Viewing $V(M)$ as the action for a zero dimensional QFT, obtain an expression for the propagator $\langle M^i_j M^k_l \rangle$ correct to lowest non-trivial order in g . [*Consider only connected diagrams.*]
- c) Now let B , C and H be further $N \times N$ matrices, each of which have zeros all along the leading diagonal. Let the elements of B and C be fermionic variables, while the entries of H are bosonic. Consider the matrix integral

$$Z(g) = \int \frac{[dM dB dC dH]}{(2\pi i)^{N(N-1)}} \exp\left(-NV(M) + i\text{tr}(HM) + \text{tr}(B[M, C])\right)$$

where the measure $[dM dB dC dH]$ indicates an integral over each entry of M , and the off-diagonal entries of B , C and H . Obtain the effective action for the eigenvalues $\{\lambda_i\}$ of M . [*Do not attempt to perform the path integral over these eigenvalues.*]

2

- a) Briefly explain the main idea of Wilson's picture of renormalization. What does it mean for an coupling to be *relevant*, *irrelevant* or *marginal*?
- b) A Euclidean scalar field theory has action $S = S_0 + S_1$, where $S_1(\phi, \Lambda)$ is a sum of polynomial interactions for ϕ and where the kinetic part of the action is

$$S_0(\phi, \Lambda) = \frac{1}{2} \int d^4x \phi(x) f\left(\frac{-\square + m^2}{\Lambda^2}\right) \phi(x)$$

with $f(z)$ any smooth function that obeys $f(z) = \Lambda^2 z$ for $z \leq 1$ and blows up rapidly when $z > 1$. Λ is a cutoff. Neglecting interactions, what is the momentum space propagator for this theory? Show that field modes whose energies are high compared to Λ do not propagate.

- c) Show that if the interaction terms obey

$$\Lambda \frac{\partial}{\partial \Lambda} e^{-S_1} = -\frac{1}{2} \int d^4p (2\pi)^4 \Lambda \frac{\partial(1/f)}{\partial \Lambda} \left[\frac{\delta^2 e^{-S_1}}{\delta \tilde{\phi}(p) \delta \tilde{\phi}(-p)} \right] \quad (\dagger)$$

then the partition function $Z(\Lambda)$ is independent of the cutoff. [*Neglect a contribution which simply rescales Z .*]

- d) By carrying out the field derivatives, interpret equation (\dagger) in terms of Feynman diagrams.

3

- a) Explain how the one-particle-irreducible electron self-energy $\Sigma(\not{p})$ is related to the exact momentum space electron propagator

$$iG(\not{p}) = \int d^D x e^{ip \cdot (x-y)} \langle \psi(x) \bar{\psi}(y) \rangle$$

in QED. What is the relation between this propagator and the physical mass of the electron?

- b) Briefly explain the difference between the *on-shell* and *minimal subtraction* (or $\overline{\text{MS}}$) renormalization schemes.
- c) In $D = 4 - \epsilon$ dimensions, the 1-loop contribution to the electron self-energy in QED is

$$\Sigma^{1\text{loop}}(\not{p}) = \frac{e^2}{8\pi^2} \int_0^1 dx (x\not{p} - 2m) \left[\frac{2}{\epsilon} + \ln \frac{4\pi e^{-\gamma_E} \mu^2}{(1-x)(m^2 - p^2 x)} \right],$$

where e is a dimensionless coupling and m is the renormalized electron mass. μ is an arbitrary renormalization scale while $\gamma_E \approx 0.577$ is the Euler-Mascheroni constant. Which counterterms are needed to absorb the divergent part of this integral?

- d) Show that in the $\overline{\text{MS}}$ renormalization scheme, the renormalized electron mass m is related to the physical electron mass m_{phys} by

$$m = m_{\text{phys}} \left[1 - \frac{e^2}{16\pi^2} \left(5 + 3 \ln \frac{\mu^2}{m_{\text{phys}}^2} \right) \right]$$

up to terms of order e^4 .

4

- a) Consider an $SU(N)$ gauge theory. Define the covariant derivative $D_\mu\psi$ acting on a field ψ in the adjoint representation of the gauge group.
- b) BRST transformations are defined as

$$\begin{aligned}\delta A_\mu &= \epsilon D_\mu c, & \delta c &= \frac{i}{2}\epsilon [c, c] \\ \delta \bar{c} &= \epsilon h, & \delta h &= 0\end{aligned}$$

where c and \bar{c} are ghost and antighost fields while h is an auxiliary (Nakanishi–Lautrup) field. ϵ is an anticommuting parameter. Show that these transformations are nilpotent when acting on an arbitrary operator $\mathcal{O}(A_\mu, c, \bar{c}, h)$ that is polynomial in the fields.

- c) Obtain the full action, including ghosts, for a pure (with no matter) Abelian gauge theory where the gauge field is constrained to obey $\partial^\mu A_\mu + A^\mu A_\mu = 0$. What are the ghost–antighost propagator and the ghost–antighost–gauge field vertex in momentum space?
- d) Explain why correlation functions of gauge invariant operators are independent of this gauge fixing condition.

END OF PAPER