

MATHEMATICAL TRIPOS      Part III

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Monday, 3 June, 2013    9:00 am to 12:00 pm

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PAPER 49

GENERAL RELATIVITY

*Attempt no more than **THREE** questions.*

*There are **FOUR** questions in total.*

*The questions carry equal weight.*

**STATIONERY REQUIREMENTS**

*Cover sheet*

*Treasury Tag*

*Script paper*

**SPECIAL REQUIREMENTS**

*None*

<p><b>You may not start to read the questions printed on the subsequent pages until instructed to do so by the Invigilator.</b></p>
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1

(a) Let  $T$  be a tensor field of type  $(r, s)$  and  $X$  a vector field. Define the Lie derivative  $\mathcal{L}_X T$ .

(b) Prove that the Lie derivative commutes with contraction and satisfies the Leibnitz rule

$$\mathcal{L}_X(S \otimes T) = (\mathcal{L}_X S) \otimes T + S \otimes (\mathcal{L}_X T)$$

where  $S$  is a tensor field of type  $(p, q)$ . [You may assume the existence of coordinates  $(t, x^1, x^2, \dots)$  such that  $X = \partial/\partial t$ ]

(c) Prove that  $\mathcal{L}_X f = X(f)$  and  $\mathcal{L}_X Y = [X, Y]$  where  $f$  is a function and  $Y$  a vector field.

(d) Let  $\omega$  be a covector field. Prove that, in any coordinate basis,

$$(\mathcal{L}_X \omega)_\mu = X^\nu \omega_{\mu, \nu} + \omega_\nu X^\nu{}_{, \mu}$$

(e) Let  $T$  be tensor field of type  $(1, 1)$ . Derive an expression for the components of  $\mathcal{L}_X T$  in an arbitrary coordinate basis. Hence or otherwise prove that

$$\mathcal{L}_X \mathcal{L}_Y T - \mathcal{L}_Y \mathcal{L}_X T = \mathcal{L}_{[X, Y]} T$$

## 2

In the linearized approximation to General Relativity it is assumed that the spacetime manifold is  $\mathbb{R}^4$  and there exist “almost inertial” coordinates  $x^\mu = (t, x^i)$  such that the metric takes the form  $g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu}$  where  $\eta_{\mu\nu} = \text{diag}(-1, 1, 1, 1)$  and  $|h_{\mu\nu}| \ll 1$ . Indices are raised with  $\eta^{\mu\nu}$  and lowered with  $\eta_{\mu\nu}$ . Let  $\bar{h}_{\mu\nu} = h_{\mu\nu} - (1/2)h\eta_{\mu\nu}$  where  $h = h^\rho{}_\rho$ . Then, in the gauge  $\partial^\nu \bar{h}_{\mu\nu} = 0$ , the linearized Einstein equation is  $\partial^\rho \partial_\rho \bar{h}_{\mu\nu} = -16\pi T_{\mu\nu}$ .

(a) Assume that matter is confined within a ball of radius  $d$  centred on the origin. Assume that the matter moves non-relativistically. Show that, for  $r = |\mathbf{x}| \gg d$

$$\bar{h}_{ij}(t, \mathbf{x}) \approx \frac{2}{r} \ddot{I}_{ij}(t - r)$$

where

$$I_{ij}(t) = \int d^3x T_{00}(t, \mathbf{x}) x^i x^j$$

(b) Consider a binary system consisting of two stars, each of mass  $M$ , in a Newtonian circular orbit of radius  $R$ . Assuming that the quadrupole formula applies to this system, calculate the average power emitted in gravitational radiation. You may approximate the energy-momentum tensor of a star at rest at the origin as  $T_{00} = M\delta^3(\mathbf{x})$ ,  $T_{0i} = T_{ij} = 0$ . [Hint: Take the positions of the stars to be  $\pm(R \cos \Omega t, R \sin \Omega t, 0)$  where  $\Omega$  is to be determined.]

(c) Explain why the binary systems which emit the most gravitational radiation are tightly bound and involve neutron stars or black holes.

## 3

(a) Consider an infinitesimal variation of the spacetime metric  $g_{ab} \rightarrow g_{ab} + \delta g_{ab}$ .

(i) Show that the corresponding change in the Levi-Civita connection is given by

$$\delta\Gamma_{bc}^a = \frac{1}{2}g^{ad}(\nabla_c\delta g_{db} + \nabla_b\delta g_{dc} - \nabla_d\delta g_{bc})$$

where  $\nabla$  is the Levi-Civita connection associated to  $g_{ab}$ .

(ii) Show that the change in the Ricci tensor is

$$\delta R_{ab} = \nabla_c\delta\Gamma_{ab}^c - \nabla_b\delta\Gamma_{ac}^c$$

[You may use the formula for the components of the Riemann tensor in a coordinate basis:

$$R^\mu{}_{\nu\rho\sigma} = \partial_\rho\Gamma_{\nu\sigma}^\mu - \partial_\sigma\Gamma_{\nu\rho}^\mu + \Gamma_{\nu\sigma}^\tau\Gamma_{\tau\rho}^\mu - \Gamma_{\nu\rho}^\tau\Gamma_{\tau\sigma}^\mu ]$$

(iii) Show that the change in the Ricci scalar is

$$\delta R = -R^{ab}\delta g_{ab} + \nabla^a\nabla^b\delta g_{ab} - \nabla^c\nabla_c(g^{ab}\delta g_{ab})$$

(b) A theory of gravity has action

$$S = \int d^4x\sqrt{-g}f(R) + S_{\text{matter}}$$

where  $f$  is a smooth function.

(i) Show that varying  $g_{ab}$  gives the equation of motion  $E_{ab} = (1/2)T_{ab}$  where  $T_{ab}$  is the energy-momentum tensor of matter and

$$E_{ab} = f'R_{ab} - f''\nabla_a\nabla_b R - f'''\nabla_a R\nabla_b R + \left(-\frac{1}{2}f + f''\nabla^c\nabla_c R + f'''\nabla^c R\nabla_c R\right)g_{ab}$$

where  $f'$  denotes  $f'(R)$  etc. [You may use without proof  $\delta g = gg^{ab}\delta g_{ab}$ ]

(ii) Explain why  $\nabla^a E_{ab} = 0$  for any metric  $g_{ab}$ . [Hint. This does not require a long calculation.]

(iii) Why does Lovelock's theorem not apply to this theory?

(iv) Determine the necessary and sufficient condition on the function  $f$  that ensures that any solution of the vacuum Einstein equation with cosmological constant  $\Lambda$  is also a vacuum solution of this theory.

4

A spacetime has metric

$$ds^2 = \frac{1}{z^2} (-f(z)^2 dt^2 + dx^2 + dy^2 + f(z)^{-2} dz^2)$$

where  $f(z) = \sqrt{1 - \alpha z^3}$  and  $\alpha$  is a constant. An orthonormal basis is defined by

$$e^0 = \frac{f}{z} dt \quad e^1 = \frac{1}{z} dx \quad e^2 = \frac{1}{z} dy \quad e^3 = \frac{1}{zf} dz$$

(a) Determine the connection 1-forms satisfying  $de^\mu = -\omega^\mu{}_\nu \wedge e^\nu$  and the curvature 2-forms defined by  $\Theta^\mu{}_\nu = d\omega^\mu{}_\nu + \omega^\mu{}_\rho \wedge \omega^\rho{}_\nu$

(b) Show that this metric satisfies the vacuum Einstein equation with a cosmological constant  $\Lambda$  whose value you should determine.

**END OF PAPER**