

MATHEMATICAL TRIPOS      Part III

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Friday 9 June, 2006    9 to 12

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PAPER 72

PHYSICAL COSMOLOGY

Attempt **THREE** questions.

There are **FOUR** questions in total.

The questions carry equal weight.

**STATIONERY REQUIREMENTS**

Cover sheet  
Treasury Tag  
Script paper

**SPECIAL REQUIREMENTS**

None

<p>You may not start to read the questions printed on the subsequent pages until instructed to do so by the Invigilator.</p>
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1 (i) The deceleration parameter is defined by

$$q_0 \equiv -\frac{\ddot{a}_0}{a_0 H_0^2},$$

where  $a_0$  is the scale factor today and  $H_0$  the Hubble constant. Prove that in a matter dominated Universe ( $\Lambda = 0$ )

$$\Omega_{m,0} = 2q_0,$$

where  $\Omega_{m,0}$  is the ratio of the matter density today to that required for a flat universe.

(ii) Hence show that:

$$\left(\frac{\dot{a}}{a_0}\right)^2 = H_0^2 \left[1 - 2q_0 + 2q_0 \frac{a_0}{a}\right].$$

(iii) Show that the solutions to the equation at (ii) are as follows:

(a)  $q_0 > 1/2$

$$H_0 \cdot t = q_0(2q_0 - 1)^{-3/2} [\theta - \sin \theta],$$

where the development angle  $\theta$  is defined by

$$1 - \cos \theta = \left(\frac{2q_0 - 1}{q_0}\right) \frac{a}{a_0}.$$

(b)  $q_0 = 1/2$

$$H_0 \cdot t = \frac{2}{3} \left(\frac{a}{a_0}\right)^{3/2}.$$

(c)  $0 < q_0 < 1/2$

$$H_0 \cdot t = q_0(1 - 2q_0)^{-3/2} [\sinh \psi - \psi],$$

where

$$\cosh \psi - 1 = \left(\frac{1 - 2q_0}{q_0}\right) \frac{a}{a_0}.$$

You may use natural units, whereby  $c = 1$ , throughout.

- 2 (i) From the relation between redshift and scale factor:

$$1 + z(t_0, t_e) \equiv \frac{a(t_0)}{a(t_e)}$$

where  $t_0$  is the time of observation and  $t_e$  is the time when a photon was emitted (or absorbed), show that the measured redshift of a distant astronomical source varies with  $t_0$  as:

$$\dot{z} = \frac{dz}{dt_0} = H(t_0)(1+z) - H(t_e)$$

where  $H(t)$  is the Hubble parameter at time  $t$ .

- (ii) Discuss, with the aid of appropriate sketches, the behaviour of the function  $\dot{z} = f(z)$  in flat cosmologies with:

- (a)  $\Omega_{m,0} = 1, \Omega_{\Lambda,0} = 0$ ;
- (b)  $\Omega_{m,0} = 0, \Omega_{\Lambda,0} = 1$ ;
- (c)  $\Omega_{m,0} = 0.3, \Omega_{\Lambda,0} = 0.7$ .

Does the redshift at which  $\dot{z} = 0$  depend on the present-day expansion rate?

- (iii) Using the relation

$$\frac{d\lambda}{\lambda} = \frac{dv}{c}$$

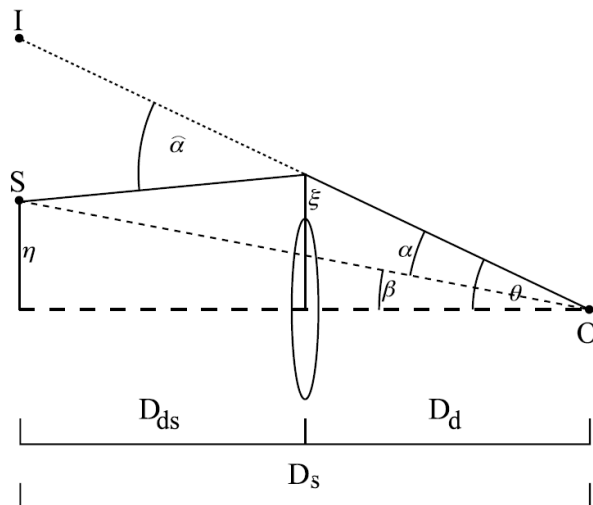
where  $d\lambda$  and  $dv$  are small intervals of wavelength and velocity, and assuming an Einstein-de Sitter cosmology ( $\Omega_{m,0} = 1; \Omega_{\Lambda,0} = \Omega_k = 0$ ), estimate the velocity shift between two spectra of the same extragalactic source at  $z = 3$  recorded at two epochs separated by 10 years. You can assume that the source has zero peculiar velocity relative to the cosmic expansion.

Note that:

- (1)  $H_0 = 70 \text{ km s}^{-1} \text{ Mpc}^{-1} \simeq 2 \times 10^{-18} \text{ s}^{-1}$ ;
- (2) 1 year  $\simeq 3 \times 10^7 \text{ s}$ .

- (iv) Comment on the feasibility of using observations of absorption lines in the Lyman alpha forest to measure directly the cosmic expansion.

**3** The figure shows a light ray propagating from the point source  $S$ , being deflected by an angle  $\hat{\alpha}$  by the gravitational potential of a lens  $L$ , and reaching the observer at  $O$ . The transverse distance of  $S$  from the optical axis connecting  $O$  and  $L$  is  $\eta$ , while  $\xi$  is the transverse distance from the optical axis of the light ray at the point of deflection. The angular separations of the source and image  $I$  from the optical axis, as seen by the observer, are  $\beta$  and  $\theta$  respectively. The angular diameter distances between the observer and the lens, the observer and the source, and the lens and the source are, respectively,  $D_d$ ,  $D_s$ , and  $D_{ds}$ .



(i) Derive the relationship between  $\hat{\alpha}$ ,  $\eta$ ,  $\xi$ ,  $D_d$ ,  $D_s$ , and  $D_{ds}$  which must be satisfied for the light ray to reach the observer at  $O$ . Similarly, derive a relationship between  $\hat{\alpha}$  and the scaled deflection angle  $\alpha$ .

(ii) Use the relationships derived at (i) to verify the one-dimensional lens equation:  $\beta = \theta - \alpha(\theta)$ . Hence, given that for a point-mass lens:

$$\hat{\alpha} = \frac{4GM}{c^2\xi}$$

use the lens equation to solve for the Einstein radius  $\theta_E$  in terms of the angular diameter distances, and for the image positions  $\theta$  in terms of  $\theta_E$  and  $\beta$ .

(iii) Given that the magnification of an image is

$$\mu = \frac{\theta}{\beta} \frac{d\theta}{d\beta},$$

derive an expression for  $\mu$  in terms of  $\theta$  and  $\theta_E$  only. Discuss the effect of gravitational lensing by a point mass as a point source is moved from (a)  $\beta = 0$ , to (b)  $0 < \beta < \theta_E$ , and (c)  $\beta \gg \theta_E$ .

(iv) Galaxies are not point-masses. A more realistic model for a lensing galaxy is given by:

$$\kappa(\theta) = \frac{\theta_E}{\theta_c} \left(1 + \frac{\theta^2}{2\theta_c^2}\right) \left(1 + \frac{\theta^2}{\theta_c^2}\right)^{-3/2}$$

where  $\kappa(\theta)$  is the convergence inside angular radius  $\theta$  with mean value:

$$\bar{\kappa}(\theta) = \frac{\theta_E}{\sqrt{\theta^2 + \theta_c^2}}$$

$\theta_c$  is the core radius and  $\theta_E$  is a constant.

- (a) Deduce  $\lim_{\theta \rightarrow \infty} \kappa(\theta)$  and  $\lim_{\theta \rightarrow 0} \kappa(\theta)$ . How does  $\lim_{\theta \rightarrow 0} \kappa(\theta)$  compare with that of a singular isothermal sphere?
- (b) What simple condition must  $\theta_c$  satisfy for multiple images to occur?
- (c) Find two sets of conditions involving  $\kappa(\theta)$  and/or  $\bar{\kappa}(\theta)$  which satisfy

$$\det \mathcal{A}(\theta) = 0,$$

where  $\mathcal{A}(\theta)$  is the Jacobian matrix of the lens mapping.

- (d) Solve one of the sets of conditions for  $\theta$  in terms of  $\theta_E$  and  $\theta_c$ .
- (e) Comment on the significance of the conditions.

4 (i) Explain what is meant by the Cosmological Principle and justify its adoption in modern cosmology. Describe three empirical observations which support it.

(ii) The term ‘Consensus Cosmology’ refers to a set of cosmological parameters which have now been determined with an accuracy of better than 10%. Among these parameters are:

- (a) the Hubble constant  $H_0 = 70 \text{ km s}^{-1} \text{ Mpc}^{-1}$ ;
- (b) the density parameters  $\Omega_{b,0} = 0.045$ ,  $\Omega_{m,0} = 0.3$ ,  $\Omega_{\Lambda,0} = 0.7$ ,  $\Omega_{k,0} = 0$ , where  $\Omega_{i,0}$  indicates the present-day fractional contribution of, respectively, baryons, matter, the cosmological constant, and curvature to the critical density  $\rho_{\text{crit}} = 3H_0^2/8\pi G$ .

Discuss the observational evidence for these values of the cosmological parameters.

(iii) From the definition of  $\rho_{\text{crit}}$  at point (ii) and the equation of state of simple fluids:

$$p_i = w_i \rho_i,$$

where  $p_i$  is the pressure and  $w_i$  is a constant for a given component of the universe, give reasons why today’s consensus cosmology may appear contrived.

**END OF PAPER**