

# **Bosonic String in Anti-de Sitter Space**

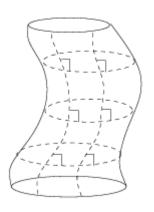
Ruoxin Bai supervised by Dr Bob Knighton

Summer Research Festival 2024,14 October 2024, Centre for Mathematical Sciences

1 <sup>st</sup> quantisation	2 <sup>nd</sup> quantisation
Path $x:I\longrightarrow \mathcal{M}$	$\begin{aligned} & \text{Field} \\ & \phi \in \Gamma(E, \mathcal{M}) \end{aligned}$
$\int \mathcal{D}x \exp(-S[x])$	$\int \mathcal{D}\phi \exp(-S[\phi])$

General Relativity is a field theory with metric  $g_{\mu\nu}$  and action  $S_{EH}$ .

It is highly non-linear and non-renormalizable. How to quantise gravity?



 $1^{st}$  quantisation + a length  $l_s$ 

We get worldsheet instead of worldline

Embedding  $X: Woldsheet \rightarrow Spacetime$ 

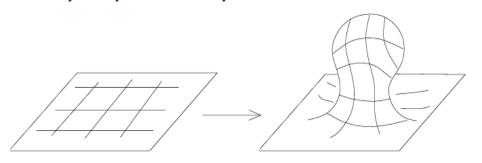


#### Polyakov action:

$${\cal S} = rac{T}{2} \int {
m d}^2 \sigma \, \sqrt{-h} \, h^{ab} g_{\mu
u}(X) \partial_a X^\mu(\sigma) \partial_b X^
u(\sigma),$$

#### Symmetries:

- Poincare symmetry in spacetime
- Reparametrization on world sheet
- Weyl symmetry



 $g_{\alpha\beta}(\sigma) \to \Omega^2(\sigma) g_{\alpha\beta}(\sigma)$ 

A very constrained system

Equations of motion and quantize! Promote *X* to an operator.

Modes satisfy the Virasoro algebra:

$$[L_n, L_m] = (n-m)L_{n+m} + \frac{c}{12}(n-1)n(n+1)\delta_{n+m,0}\mathbf{1}$$

- 2D field theory are well-understood and rich.
- D=26
- massless spin 2 particle is Einstein gravity. Reduce to GR.

### What is Anti-de Sitter Space?

The anti-de Sitter space of signature (p, q) can then be isometrically embedded in the space  $\mathbb{R}^{p,q+1}$  with coordinates  $(x_1, ..., x_p, t_1, ..., t_{q+1})$  and the metric

$$ds^2 = \sum_{i=1}^p dx_i^2 - \sum_{j=1}^{q+1} dt_j^2$$

as the quasi-sphere

$$\sum_{i=1}^p x_i^2 - \sum_{j=1}^{q+1} t_j^2 = -lpha^2,$$

where lpha is a nonzero constant with dimensions of length (the radius of curvature). This is a

Lorentzian manifold with constant negative curvature. Universe with negative cosmological constant.



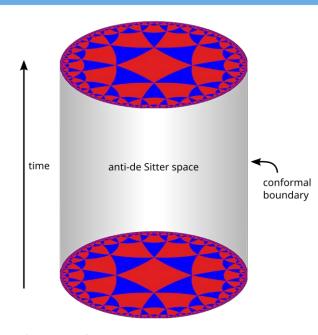
### 3-dimensional Anti-de Sitter Space

3 dimensional anti-de Sitter space:

$$X_{-1}^2 + X_0^2 - X_1^2 - X_2^2 = 1.$$

Consider Matrices paramterised by:

$$g = \begin{pmatrix} X_{-1} + X_1 & X_0 - X_2 \\ -X_0 - X_2 & X_{-1} - X_1 \end{pmatrix},$$



Set determinant to 1, the hyperboloid is  $SL(2,\mathbb{R})$ 

$$ds^2 = -\cosh^2 \rho dt^2 + d\rho^2 + \sinh^2 \rho d\phi^2.$$

If the string propagates in a Lie group,  $X:WS \to G$ . Think G as a symmetry, we have conserved currents.

$$[J_m^a, J_n^a] = k \delta^{ab} \delta_{m+n,0} + i f_c^{ab} J_{m+n}^c$$

Kac-Moody algebra.

States fall in representations  $\mathfrak g$  and treat currents as oscillators.

Wess-Zumino-Witten model.

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Why? AdS/CFT correspondence.

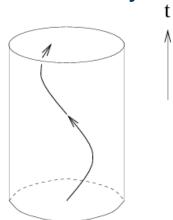


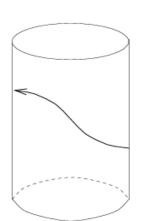
WZW action:

$$S = \frac{k}{8\pi\alpha'} \int d^2\sigma \operatorname{Tr}\left(g^{-1}\partial g g^{-1}\partial g\right) + k\Gamma_{WZ}$$

where the last term is the Wess-Zumino term.

Consider symmetries and classical geodesics:



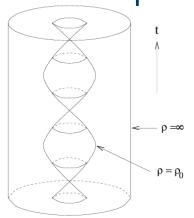


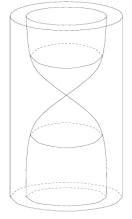
What about strings? Spectral flow:

Given a trajectory  $g(\sigma, \tau)$ , we can define new trajectories by the operation:

$$\sigma^{w}(g) = e^{-w\sigma^{+}t^{3}}g(\sigma,\tau)e^{-w\sigma^{-}t^{3}}$$

This corresponds to:  $t \rightarrow t + w\tau$   $\phi \rightarrow \phi + w\sigma$ 





$$w \in Z$$

Stretches along t, winds along  $\phi$ 



# Thank you for listening!

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